



# Photoelectrochemical water oxidation over $\text{TiO}_2$ nanotubes modified with $\text{MoS}_2$ and $\text{g-C}_3\text{N}_4$

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## Full Research Paper

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## Abstract

$\text{TiO}_2$  nanotube arrays (TNAs) have been studied for photoelectrochemical (PEC) water splitting. However, there are two major barriers of TNAs, including a low photo-response and the fast charge carrier recombination in TNAs, leading to poor photocatalytic efficiency. Through a comparison of  $\text{MoS}_2/\text{TNAs}$  and  $\text{g-C}_3\text{N}_4/\text{TNAs}$ , it was found that TNAs modified with  $\text{MoS}_2$  and  $\text{g-C}_3\text{N}_4$  exhibited a current density of, respectively, 210.6 and 139.6  $\mu\text{A}\cdot\text{cm}^{-2}$  at an overpotential of 1.23 V vs RHE, which is 18.2 and 12 times higher than that of pure TNAs under the same conditions. The stability of the  $\text{MoS}_2/\text{TNAs}$  heterojunction is higher than that of  $\text{g-C}_3\text{N}_4/\text{TNAs}$ .

## Introduction

Hydrogen energy has become a target pursued in the energy development strategies of many countries and regions. Hydrogen is often synthesized via hydrocarbon compounds or water electrolysis [1]. Methods to produce hydrogen via electrochemical or photo-electrochemical (PEC) water splitting are considered a future direction of renewable fuel development [2–4]. The use of solar energy to activate catalytic materials to separate water for creating clean fuels has been developed for about a decade [5,6]. Water splitting is carried out in solutions rich in  $\text{H}^+$  ions

to the conduct hydrogen evolution reaction (HER) process or in rich  $\text{OH}^-$  solutions for the oxygen evolution reaction (OER) process [7–9]. However, the electrode material must be extremely durable and nearly chemically inert to be able to withstand highly acidic or basic environments. Therefore, noble metals such as Pt, Pd, Au and Ag with suitable chemical properties, such as inertness, good resistance against corrosion and good electrical conductivity have been widely used in water splitting reactions [10,11]. However, noble metals are still rare

and expensive materials, and their application as electrode materials is considered to be not optimal [10]. Therefore, the study of a materials with high-performance in PEC water splitting, which could replace noble metals are a research interest.

Photocatalytic technology uses semiconductors for effective approaches to the degradation of dyes and antibiotics, the removal of pollutant gases, and water splitting to produce hydrogen using solar energy [12–17]. Among such semiconductors, TiO<sub>2</sub> nanotube arrays (TNAs) of 2–100 nm in diameter and 1–2 μm in length, are often used for efficient PEC applications exploiting advantages such as chemical stability, less toxicity and suitable cost [18–21]. However, there are two disadvantages affecting directly their photocatalytic ability. (i) TNAs only respond to ultraviolet (UV) light [22–24], and (ii) they exhibit fast carrier recombination [25]. Recently, the development of new heterojunction architectures through coupling TNAs with other semiconductor materials, especially low-bandgap semiconductors, led to a reduction of the required amounts of noble metals and materials such as CdS or ZnS [26–29]. There are many low-bandgap semiconductors that were coupled with TNAs, including MoS<sub>2</sub>, WS<sub>2</sub>, MoSe<sub>2</sub>, g-C<sub>3</sub>N<sub>4</sub>, Cu<sub>2</sub>O, and CuO. MoS<sub>2</sub> is a semiconductor with a narrow bandgap (1.9 eV at room temperature) exhibiting unique chemical, thermal, and charge transport properties, which can shift the light absorption of TiO<sub>2</sub> to the visible region [29–32]. An emerging new material in optoelectronics is g-C<sub>3</sub>N<sub>4</sub> (bandgap of 2.65–2.7 eV) because it has an appropriate band structure with suitable energy levels regarding TiO<sub>2</sub>, which can improve the charge transfer states [33,34]. These two low-bandgap semiconductors improved considerably the PEC water splitting efficiency [35,36]. However, the fabrication of MoS<sub>2</sub>/TNAs and g-C<sub>3</sub>N<sub>4</sub>/TNAs has many disadvantages such as high synthesis temperatures, the requirement of a binder, or expensive synthesis equipment [29,36–38].

In this study, we compare properties and PEC water splitting efficiency of TNAs combined with the typical 2D materials MoS<sub>2</sub> and g-C<sub>3</sub>N<sub>4</sub> obtained with the same synthesis procedure. Insightful studies about optical and electronic properties have been conducted to explain clearly the difference between these composite materials

## Experimental

### Materials and chemicals

Chemicals and materials for the synthesis and characterization include Ti foil (1 cm × 2 cm), hydrochloric acid (HCl), sodium hydroxide (NaOH), DI water, acetone ((CH<sub>3</sub>)<sub>2</sub>CO), ethanol (C<sub>2</sub>H<sub>5</sub>OH), ammonium fluoride (NH<sub>4</sub>F), *N*-acetyl-L-cysteine, ammonium heptamolybdate ((NH<sub>4</sub>)<sub>6</sub>Mo<sub>7</sub>O<sub>24</sub>), thiourea (CH<sub>4</sub>N<sub>2</sub>S), nitrogen gas, melamine, and nafion solution. All

chemicals and materials were purified and used without further treatment.

### Preparation of materials

The individual materials including TNAs, MoS<sub>2</sub>, and g-C<sub>3</sub>N<sub>4</sub> were synthesized as described earlier [39–41]. To combine with TNAs, 5 mg of MoS<sub>2</sub> or g-C<sub>3</sub>N<sub>4</sub> powder was dispersed in 2 mL of a solution containing 50 vol % ethanol and 50 vol % nafion solution as described in [40]. The solution was stirred for 30 min before ultrasonic treatment for 3 h to obtain a homogeneous solution. Next, 0.2 mL of either of these solutions was used to coat the surface of TNAs via spin coating. The samples were denoted as MoS<sub>2</sub>/TNAs and g-C<sub>3</sub>N<sub>4</sub>/TNAs. Then, the samples were annealed in nitrogen gas at 60 °C for 12 h to obtain a stable electrode for the investigation processes.

### Characterization of materials

The morphology, the phase, and the vibrational characteristics of the surface functional groups of the materials were observed by field-emission scanning electron microscopy (FESEM), X-ray diffraction (XRD), and Fourier-transform infrared spectroscopy (FTIR). Diffuse reflectance spectroscopy (DRS) was carried out to measure the optical bandgap of the semiconductor materials through the Tauc method using the absorption coefficient  $\alpha$  of the material, according to Equation 1 [42]:

$$(\alpha h\nu)^{\frac{1}{\gamma}} = B(h\nu - E_g), \quad (1)$$

where  $h$ ,  $\nu$ ,  $E_g$ , and  $B$  are the Planck constant, the frequency of the photon, the bandgap energy, and a constant, respectively;  $\gamma = 1/2$  for materials with direct bandgap and  $\gamma = 2$  for semiconductor materials with an indirect bandgap.

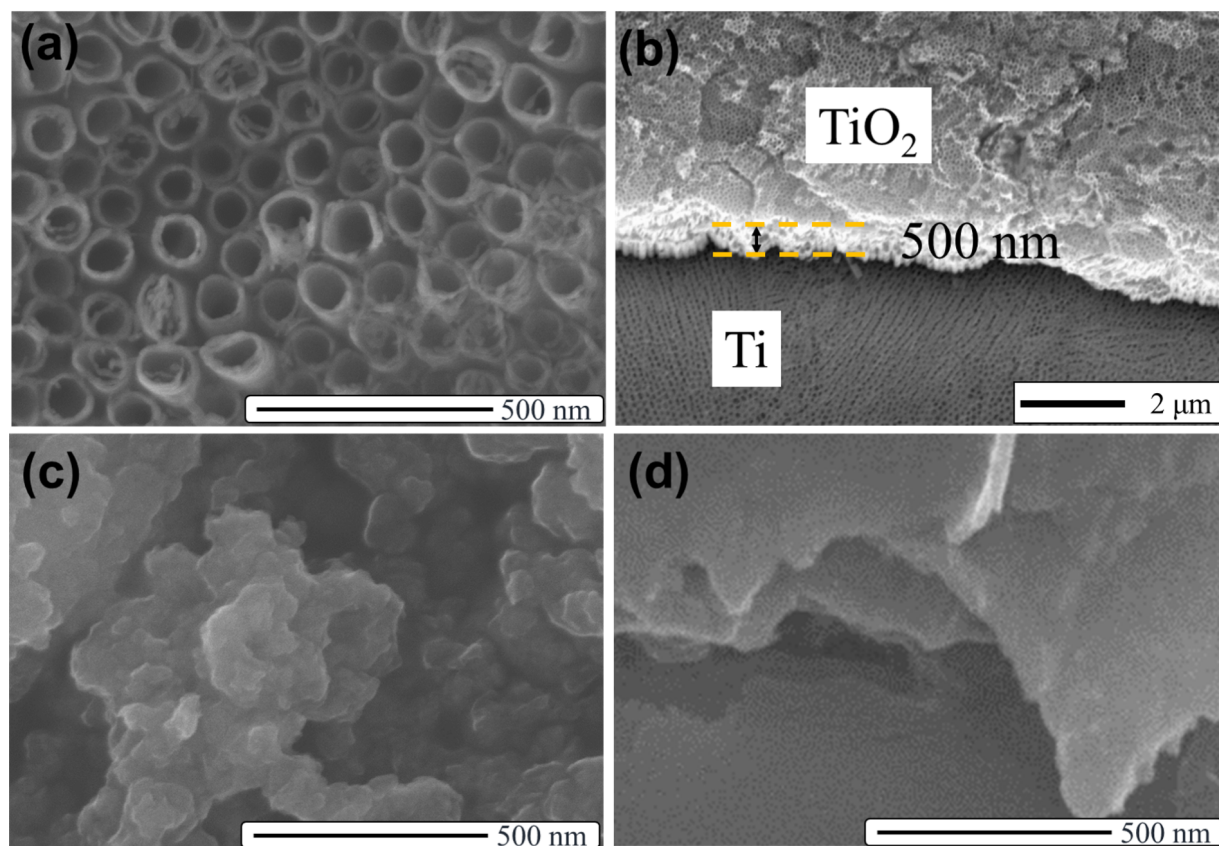
### PEC activity evaluation

The PEC water splitting efficiency of the materials was evaluated through a three-electrode PEC cell using a Biologic SP-200 potentiostat. The three electrodes were a Pt counter electrode, a Ag/AgCl 3 M reference electrode, and a MoS<sub>2</sub>/TNAs or g-C<sub>3</sub>N<sub>4</sub>/TNAs working electrode in a 1 M Na<sub>2</sub>SO<sub>4</sub> (pH 7.4) electrolyte solution. The light source used in this study was a 150 W Xe lamp (ABET Instruments) with a calibrated luminous intensity of 100 mW·cm<sup>−2</sup> and a UV filter cutting at 380 nm.

## Results and Discussion

### Characterizations of materials

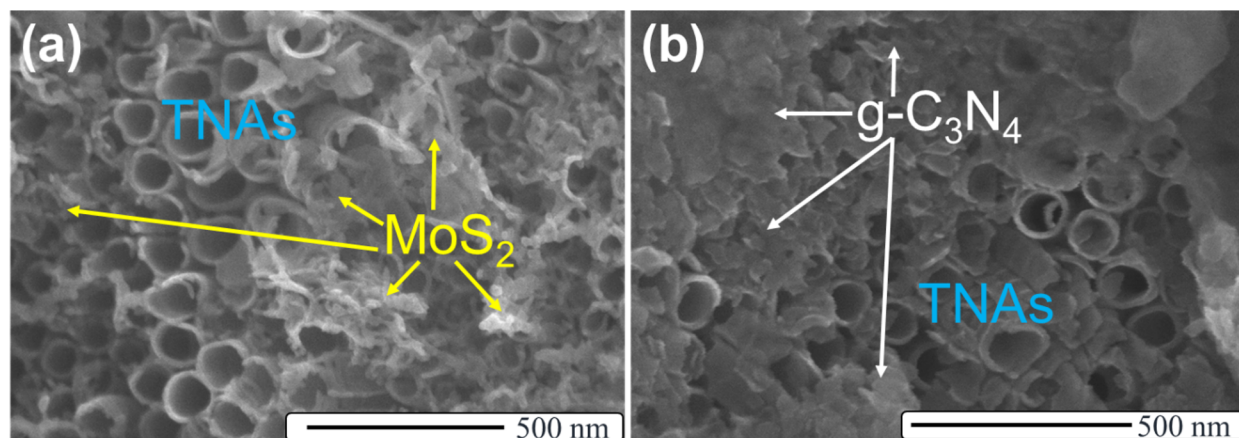
Figure 1a displays the morphology of TNAs, which have a uniform distribution of nanotubes with average diameters ranging from 80–100 nm and a length of 500 nm (Figure 1b). The MoS<sub>2</sub> material exhibits the stacked layers of 2D materials



**Figure 1:** SEM images of TNAs (a, b),  $\text{MoS}_2$  (c), and  $\text{g-C}_3\text{N}_4$  (d).

(Figure 1b). This agrees with the results of previous publications in which hydrothermal methods were applied [24–26]. The SEM image of the  $\text{g-C}_3\text{N}_4$  material shows the uniform nano-sheets that were fabricated by the melamine pyrolysis method (Figure 1c).

After the deposition of 2D materials  $\text{MoS}_2$  and  $\text{g-C}_3\text{N}_4$  onto the TNAs substrate, we examined the morphology of these hetero-structures by using SEM (Figure 2). There are some small pieces that are randomly distributed on the surface of TNAs in Figure 2a, which were attributed to be  $\text{MoS}_2$ . There is a similar



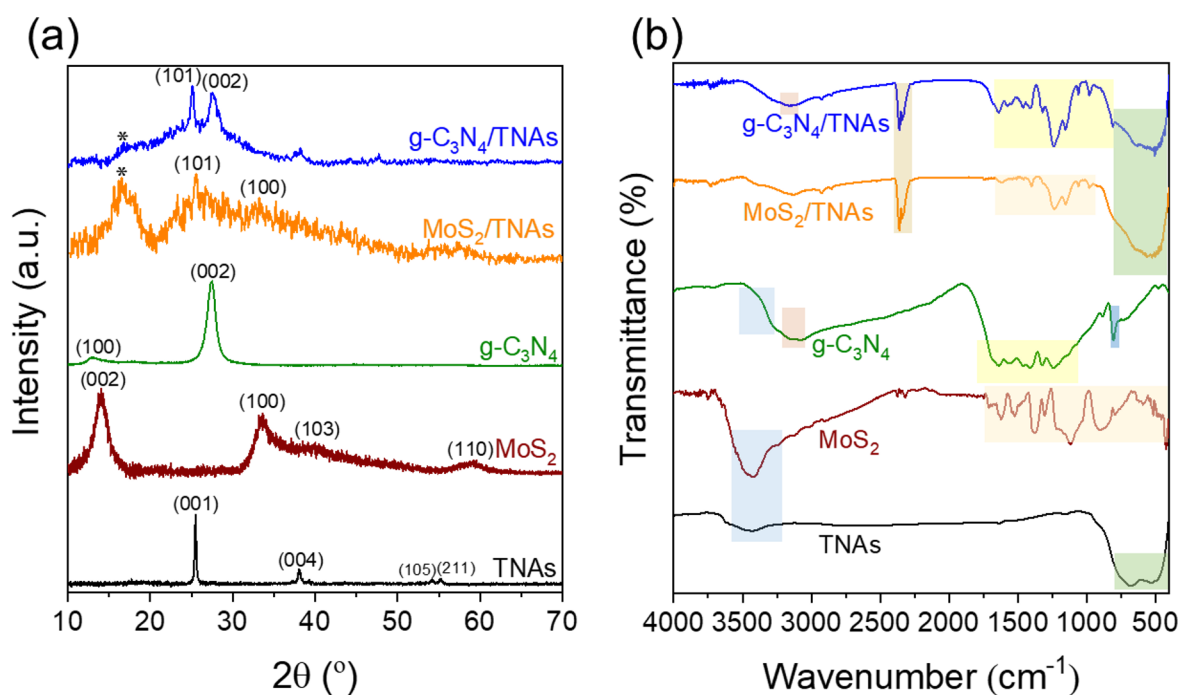
**Figure 2:** SEM images of  $\text{MoS}_2/\text{TNAs}$  (a), and  $\text{g-C}_3\text{N}_4/\text{TNAs}$  (b).

result in the SEM image of g-C<sub>3</sub>N<sub>4</sub>/TNAs in Figure 2b. However, the distribution of g-C<sub>3</sub>N<sub>4</sub> on the surface of the TNAs is denser than that of MoS<sub>2</sub>. Besides, the morphology of the TNAs did not change. The appearance of MoS<sub>2</sub> and g-C<sub>3</sub>N<sub>4</sub> has also been confirmed by EDS measurements and element mapping (Figure S1 and Figure S2, Supporting Information File 1).

Figure 3a shows that diffraction peaks of TNAs at  $2\theta = 25.45^\circ$ ,  $37.07^\circ$ ,  $39.24^\circ$ ,  $54.10^\circ$ , and  $55.25^\circ$ , attributed to the (101), (103), (004), (112), (105), and (211) planes of TiO<sub>2</sub>, respectively [JCPDS No. 21-1272]. Besides, the XRD pattern of MoS<sub>2</sub> exhibits diffraction peaks at  $13.97^\circ$ ,  $33.56^\circ$ ,  $40.24^\circ$ , and  $59.25^\circ$ , corresponding to the (002), (100), (103), and (110) planes, respectively, of the 2H phase of MoS<sub>2</sub> [JCPDS No. 37-1492]. The pristine g-C<sub>3</sub>N<sub>4</sub> shows two distinct characteristic peaks at  $2\theta = 12.9^\circ$  and  $27.45^\circ$ , assigned to the (100) and (002) planes, respectively [43,44]. The XRD diffraction results show the simultaneous appearance of diffraction peaks at  $2\theta = 25.45^\circ$ , typical for the (101) planes of TNAs, and at  $33.56^\circ$  for the (001) planes of MoS<sub>2</sub>. Besides, the diffraction peak at  $2\theta = 16.45^\circ$  characterizes the semi-crystalline structure of perfluorocarbon chains from nafion films [45]. Notably, the diffraction peak of the MoS<sub>2</sub> material at  $2\theta = 13.97^\circ$ , which is typical for the (002) lattice plane, is lost after the formation of the MoS<sub>2</sub>/TNAs heterostructure. This could be explained by the very small amount of MoS<sub>2</sub> (5 mg) loaded onto the TNAs. Therefore, it will be difficult to identify the MoS<sub>2</sub> phase in the composite

from the XRD pattern of MoS<sub>2</sub>/TNAs. The (002) plane indicates the multilayer structure of MoS<sub>2</sub> materials, the (001) plane indicates a monolayer structure of MoS<sub>2</sub> [37,46]. Therefore, the disappearing (002) reflection and the remaining (001) reflection show that the ultrasonic treatment peeled the MoS<sub>2</sub> material into thinner layered structures. This is in agreement with the SEM images, in which material with rather small and thinner structures scattered on the surface of TNAs was observed.

The functional groups and chemical bonds of the as-prepared materials were determined by using FTIR spectroscopy, as shown in Figure 3b. The formation of TiO<sub>2</sub> on the Ti foil is indicated by the vibrations of the Ti–O bond in the wavenumber region from 450 to 750 cm<sup>−1</sup> [47]. The bonding characteristics in the MoS<sub>2</sub> material are presented by Mo–S vibration peaks between 1620 and 420 cm<sup>−1</sup> [48]. Also, FTIR spectroscopy is used as an extremely effective technique for studying the vibrational states of organic bonds in g-C<sub>3</sub>N<sub>4</sub> materials using the vibrational peaks of C–N bonds between 1640 and 1200 cm<sup>−1</sup>. A very strong absorption peak at 807.2 cm<sup>−1</sup> characterizes the fluctuation of tri-s-triazine of g-C<sub>3</sub>N<sub>4</sub>. Vibrational peaks in the 3200 cm<sup>−1</sup> region attributed to fluctuations of the C–N group also appeared [49]. Figure 3b shows the bonding states in the MoS<sub>2</sub>/TNAs and g-C<sub>3</sub>N<sub>4</sub>/TNAs heterostructures. The results show that, in addition to the typical bonding of the TNAs substrate such as Ti–O bonds, there are vibrations of composites of TNAs with MoS<sub>2</sub> (between 420 and 1620 cm<sup>−1</sup>) and g-C<sub>3</sub>N<sub>4</sub>



**Figure 3:** XRD pattern (a) and FTIR spectra (b) of as-synthesized samples.

(between 1200 and 1640  $\text{cm}^{-1}$  for C–N bonds and 807  $\text{cm}^{-1}$  for tri-s-triazine subunit). The peaks in the wavenumber range between 3400 and 1625  $\text{cm}^{-1}$  of all samples are typical for stretching vibrations of the O–H bonds, which are caused by air humidity.

To confirm the ability of the heterojunctions to enhance absorption in the visible-light region, the DRS spectra and Tauc plots were recorded and are presented in Figure 4. It can be easily observed in Figure 4a that the TNAs show a strong absorption edge at 393 nm. This means that TNAs are only activated by near-UV irradiation. In contrast, the g-C<sub>3</sub>N<sub>4</sub> sample shows an absorption edge at 464 nm. Meanwhile, MoS<sub>2</sub> exhibits strong absorption from the UV region extending to the entire visible-light region. It can be seen that the loading of both MoS<sub>2</sub> and g-C<sub>3</sub>N<sub>4</sub> on the TNAs surface extended the absorption to the visible-light range. Specifically, the absorption edge of the g-C<sub>3</sub>N<sub>4</sub>/TNAs and MoS<sub>2</sub>/TNAs samples shifted to 442 and 425 nm, respectively. Besides, a remarkable feature of the DRS spectrum of MoS<sub>2</sub> is a superior absorption ability in the whole visible-light range in comparison to that of the remaining samples. To evaluate the optical bandgap energy of TNAs and g-C<sub>3</sub>N<sub>4</sub>, Tauc plots were extrapolated in Figure 4b. The bandgap values of TNAs, g-C<sub>3</sub>N<sub>4</sub>, and MoS<sub>2</sub> were calculated as about 3.15, 2.67, and 1.47 eV, respectively. These results are agreement with previous publications [50–52].

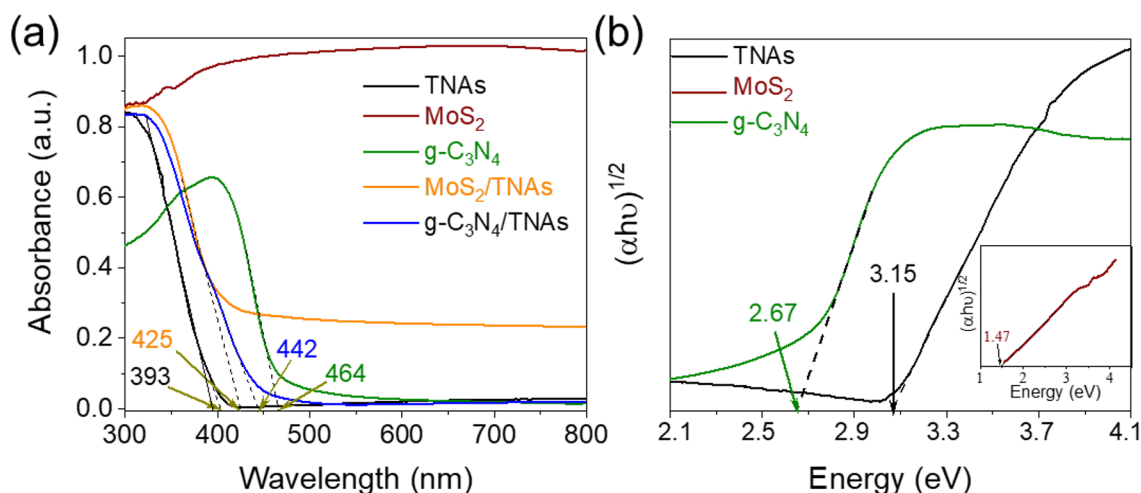
Figure 5 shows the results of electrochemical impedance spectroscopy (EIS), that is, Nyquist and Mott–Schottky plots of the materials, which give information about the charge transfer mechanism at the interface. In Figure 5a, the Nyquist plots of the samples all exhibit only single semicircular shape, which shows the charge transfer resistance equivalent to the polariza-

tion resistance. This result also demonstrates a unique interaction of the electrode surface and the electrolyte solution. Furthermore, the g-C<sub>3</sub>N<sub>4</sub> sample shows the semicircle with the largest radius, followed by TNAs and MoS<sub>2</sub>, which indicates the low interaction of these materials with the electrolyte. However, after coupling, the g-C<sub>3</sub>N<sub>4</sub>/TNAs sample shows a semicircle with smaller radius compared than that of g-C<sub>3</sub>N<sub>4</sub> or TNAs. The Nyquist plot of the MoS<sub>2</sub>/TNAs sample shows the smallest semicircle radius of all samples. This result indicates an increase in carrier density during the reaction with the electrolyte solution, leading to a decrease in resistance of the g-C<sub>3</sub>N<sub>4</sub>/TNAs and MoS<sub>2</sub>/TNAs samples. This could be explained by the enhanced optical interaction, causing an increase of the carrier concentration in MoS<sub>2</sub>/TNAs sample such in Figure 4.

The mechanism for the enhanced activity of the heterojunctions can be explained by the Mott–Schottky results in Figure 5b,c. Generally, all samples show positive slopes, which proves that they are n-type semiconductors [53]. Equation 2 shows the Mott–Schottky relationship involving the apparent capacitance as a function of the potential under depletion conditions [54]:

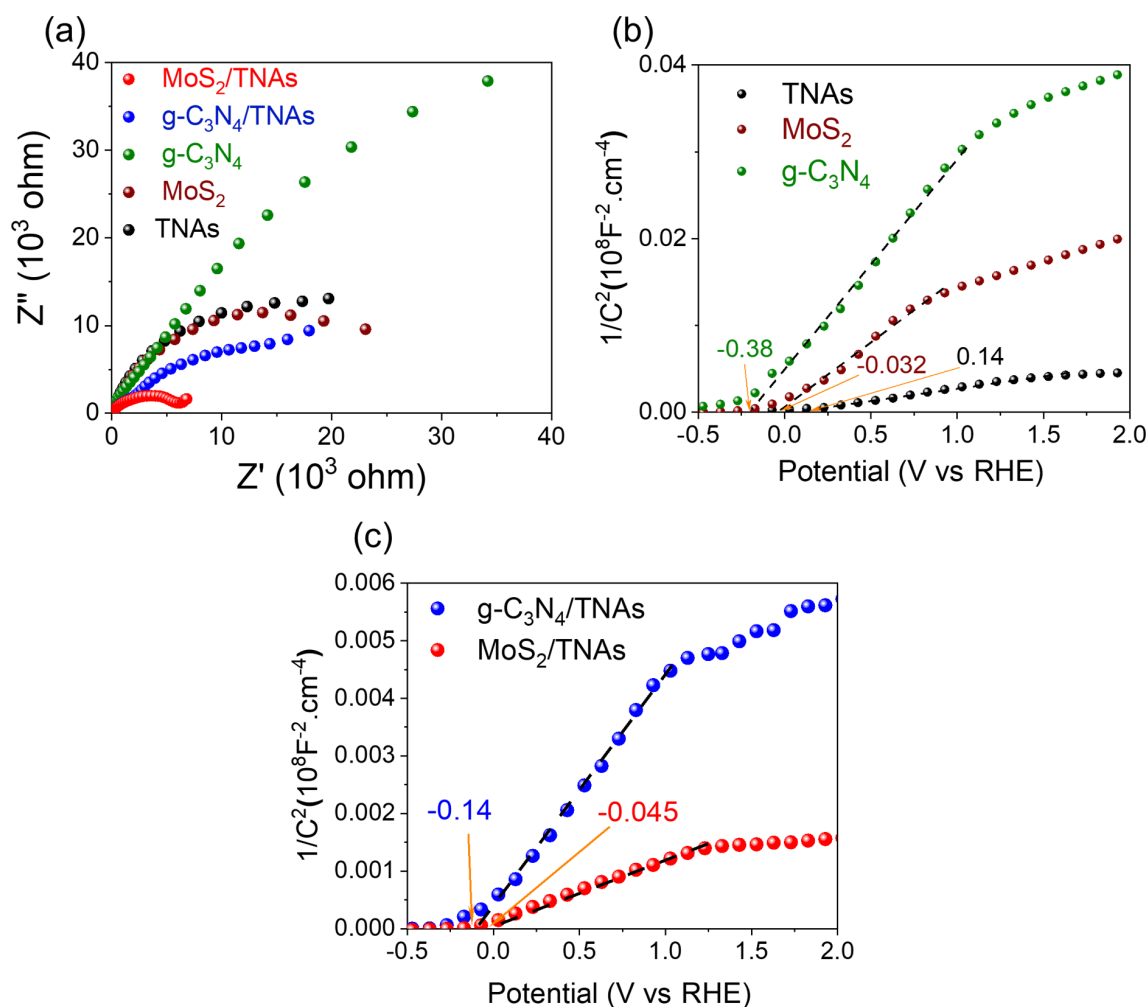
$$C^{-2} = \frac{2}{e\epsilon\epsilon_0 N_d A^2} \left( V_a - V_{fb} - \frac{kT}{e} \right), \quad (2)$$

where  $C$ ,  $\epsilon$ ,  $\epsilon_0$ ,  $N$ ,  $A$ ,  $V_a$ ,  $V_{fb}$ ,  $k$ , and  $T$  are the capacitance of the space charge region, the dielectric constant of the semiconductor, the vacuum permittivity, the donor density, the area of interface or the electrode, the applied and flat band potentials, the Boltzmann constant, and the temperature, respectively.



**Figure 4:** Comparison of the optical properties of as-synthesized materials through DRS spectra (a) and Tauc plots (b).





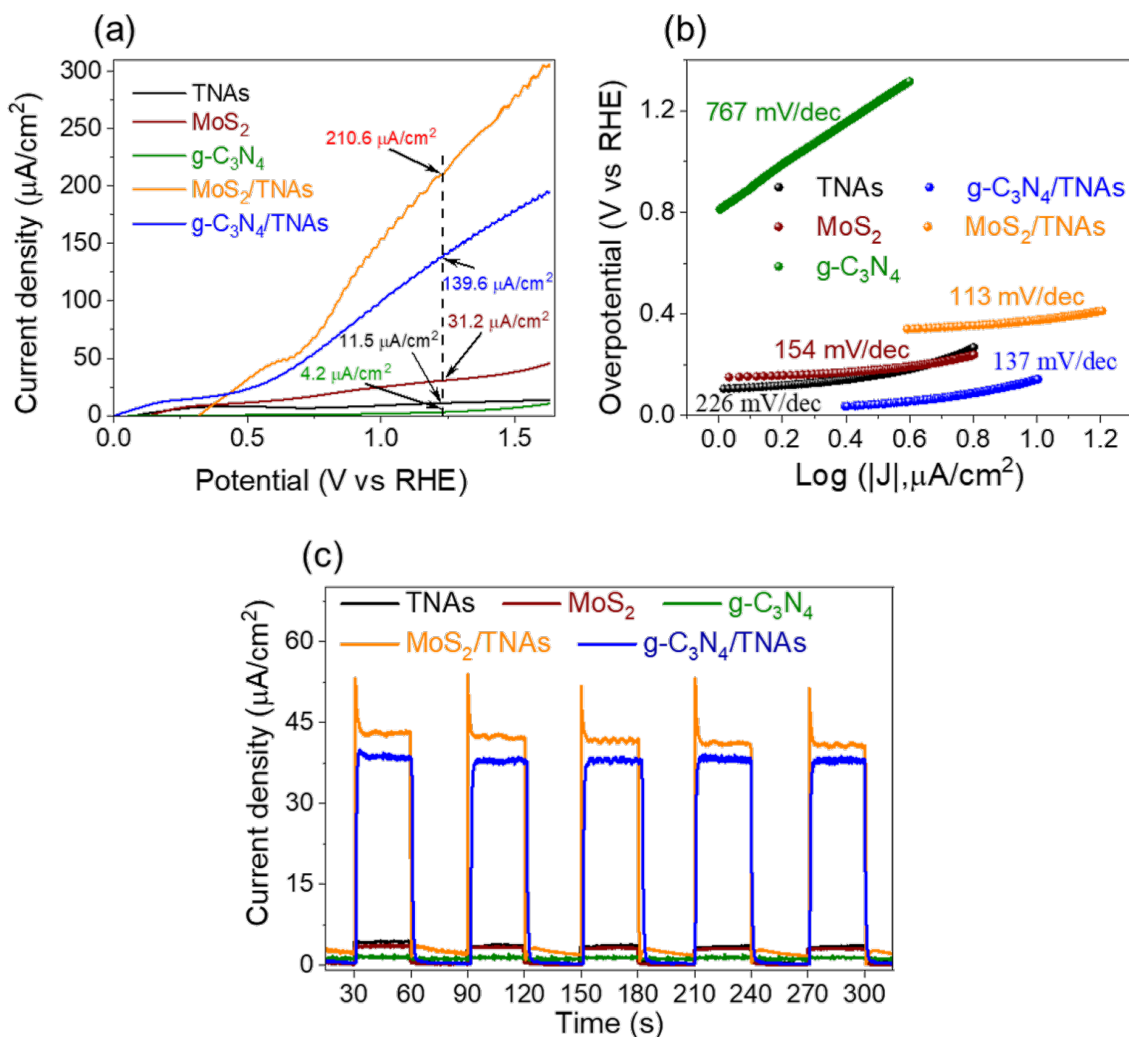
**Figure 5:** EIS spectra (a), Mott–Schottky plots of pristine materials (b) and heterostructures (c).

The plot of  $1/C^2$  vs  $V$  shows an intercept of the x-axis, which corresponds to the flat band potential ( $E_{fb}$ ), that is, the conduction band maximum (CBM) level of the material. The Mott–Schottky plots of TNAs,  $g-C_3N_4$ , and  $MoS_2$  samples yield  $E_{fb}$  (or CBM) values of 0.14,  $-0.38$ , and  $-0.032$  V vs RHE, respectively. It can be found that the  $E_{fb}$  values of  $g-C_3N_4$  and  $MoS_2$  are significantly more negative than that of TNAs, which can facilitate the migration of electrons from  $g-C_3N_4$  and  $MoS_2$  to TNAs. Furthermore, the  $E_{fb}$  values of  $g-C_3N_4$ /TNAs and  $MoS_2$ /TNAs are shifted to  $-0.14$  and  $-0.045$  V vs RHE (Figure 5c). The heterostructures express much more negative  $E_{fb}$  values than pristine TNAs, which is attributed to the enhanced electron density, the depletion of the  $E_{fb}$ , and electron–hole recombination [55].

### PEC characterizations of materials

Figure 6 shows the linear sweep voltammetry (LSV) curves, Tafel slopes, and the photo-response of the samples. Figure 6a

shows that the current density of all materials is grows linearly with the applied potential under visible-light excitation. The onset potentials of the of TNAs,  $g-C_3N_4$ , and  $MoS_2$  are 0.08, 0.16, and 0.14 V vs RHE, respectively. Further, for the OER (1.23 V vs RHE), the current densities of TNAs,  $g-C_3N_4$ , and  $MoS_2$  are 11.5, 4.2, and 31.2  $\mu A/cm^2$ , respectively. The onset potential values of  $g-C_3N_4$ /TNAs and  $MoS_2$ /TNAs are significantly shifted to  $-0.76$  and 0.008 V, respectively. In addition, the current density also exhibited an improvement with values of 139.6 and 210.6  $\mu A/cm^2$  at 1.23 V for  $g-C_3N_4$ /TNAs and  $MoS_2$ /TNAs, respectively, which shows their superiority in the PEC water oxidation reaction. The LSV results are also consistent with the previous results from EIS analysis and the Mott–Schottky results (Figure 5). The PEC activity of  $MoS_2$ /TNAs in this study is higher than that of  $MoS_2$ /TNAs synthesized by using a PVA binder agent in [36]. However, the direct combination of  $g-C_3N_4$  with TNAs at a relatively high fabrication temperature (500 °C for 2 h) in [35] yielded better results



**Figure 6:** LSV plots (a), Tafel slopes (b), and photo-response (c) of the materials.

better than those of this study. The investigation of the stability of the PEC electrodes from  $\text{MoS}_2/\text{TNAs}$  and  $\text{g-C}_3\text{N}_4/\text{TNAs}$  is described in Figure S3, Supporting Information File 1. After every PEC test cycle, we immersed the PEC electrode in DI water for 1 h and let it dry completely before the next test. We can conclude that the stability of the  $\text{MoS}_2/\text{TNAs}$  heterojunction is higher than that of the  $\text{g-C}_3\text{N}_4/\text{TNAs}$  heterojunction. The decrease in catalytic activity of the PEC electrodes is explained by the leaching of the catalysts  $\text{MoS}_2$  and  $\text{g-C}_3\text{N}_4$  after each activity measurement.

The Tafel slope is considered as an important parameter to evaluate the PEC activity in water splitting. A smaller Tafel slope value indicates a more rapid increase in the reaction rate of the electrode. Figure 6b shows the high Tafel slope values of the individual materials, TNAs,  $\text{g-C}_3\text{N}_4$ , and  $\text{MoS}_2$ , of 226, 767, and 154 mV/dec, respectively. After the modification, the Tafel

slope value of  $\text{g-C}_3\text{N}_4/\text{TNAs}$  is only about 137 mV/dec, while the best Tafel slope value of  $\text{MoS}_2/\text{TNAs}$  is only 113 mV/dec. Furthermore, the photocurrent of the materials was evaluated through the assessment of the photo-response under visible-light irradiation at 0.63 V in Figure 6c. A current density of about 38.6  $\mu\text{A}/\text{cm}^2$  was obtained with the  $\text{g-C}_3\text{N}_4/\text{TNAs}$  even after five cycles, which is nearly ten times higher than that of pure TNAs. The current density of  $\text{MoS}_2/\text{TNAs}$  is even higher than that of  $\text{g-C}_3\text{N}_4/\text{TNAs}$  reaching 43.4  $\mu\text{A}/\text{cm}^2$  after five cycles. These results indicate that enhancement of the optical interaction in  $\text{MoS}_2/\text{TNAs}$  heterostructures is stronger than that in  $\text{g-C}_3\text{N}_4/\text{TNAs}$  [56,57]. Further, the current density increases sharply and decreases rapidly within a few seconds for  $\text{MoS}_2/\text{TNAs}$  under light, which can be explained as follows: The photocurrent density of  $\text{MoS}_2/\text{TNAs}$  promptly increased because of the efficient separation of the  $\text{e}^-$ - $\text{h}^+$  pairs at the interfaces between TNAs and  $\text{MoS}_2$  [58] and the rapid transfer

of the photo-induced electrons from MoS<sub>2</sub> to the TNAs electrode [59]. This result is in agreement with the EIS results in Figure 5a, where the arc radius of the Nyquist plot of MoS<sub>2</sub>/TNAs was the smallest, indicating that MoS<sub>2</sub>/TNAs effectively decreased the resistance of the TNAs and, thus, speeded up the charge transfer on the photoelectrode. These arguments are consistent with results previously published in [58].

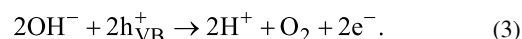
Figure 7 presents the energy band diagram structure of the MoS<sub>2</sub>/TNAs and g-C<sub>3</sub>N<sub>4</sub>/TNAs heterojunctions based on the DRS and Mott–Schottky analysis results, which are summarized in Table 1.

**Table 1:** The CBM and  $E_g$  values of the materials.

Sample	CBM level (V vs RHE, pH 7.3)	$E_g$ (eV)
TNAs	0.14	2.99
MoS <sub>2</sub>	−0.032	1.47
g-C <sub>3</sub> N <sub>4</sub>	−0.38	2.63

It is easily observed from Figure 7 that the heterostructures formed upon incorporation of TNAs with MoS<sub>2</sub> and g-C<sub>3</sub>N<sub>4</sub> are all of type II. Type-II heterostructures promote the migration of  $h^+$  and  $e^-$  under visible-light irradiation. Electrons can move from the conduction band (CB) of MoS<sub>2</sub> or g-C<sub>3</sub>N<sub>4</sub> to the CB of TNAs in MoS<sub>2</sub>/TNAs or g-C<sub>3</sub>N<sub>4</sub>/TNAs, respectively. In contrast, holes will migrate from the valence band (VB) of TNAs to the VB of MoS<sub>2</sub> or g-C<sub>3</sub>N<sub>4</sub>. Therefore, the recombination of the photo-generated  $e^-$ – $h^+$  pairs is reduced. In this contribution, the PEC water splitting reactions take place in a neutral media, which is well known to occur via two processes, including the oxidation and reduction reactions at, respectively, the anode and cathode described by Equation 3 and Equation 4.

Oxidation reaction at the anode:



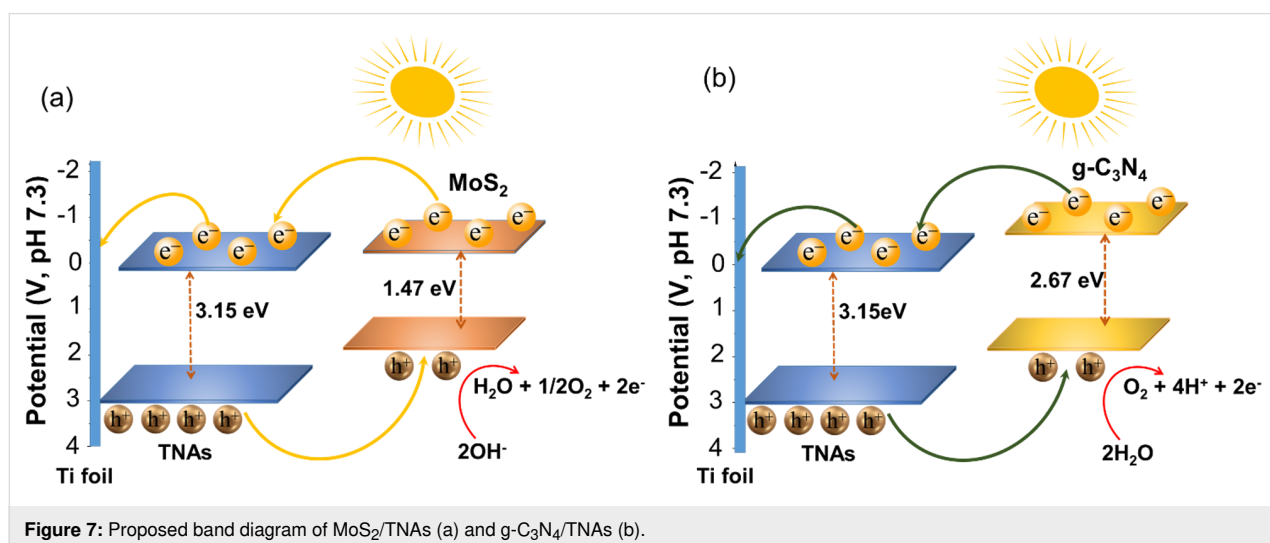
Reduction reaction at the cathode:



Carrying out the reactions in a neutral medium also contributes to the increased durability of the electrodes. However, the lack of initial  $\text{H}^+$  concentration can reduce the efficiency of the  $\text{H}_2$  production. For an effective water splitting process, the oxidation reaction of  $\text{OH}^-$  ions in the electrolyte needs to take place at the anode to generate  $e^-$  and  $\text{H}^+$  ions along with  $\text{O}_2$ . The  $e^-$  current will immediately migrate to the cathode to perform reduction reactions. At that time,  $\text{H}^+$  will also be reduced at the cathode to form  $\text{H}_2$ . The higher the efficiency of the oxidation reaction, the more  $e^-$  are generated and the stronger the  $\text{H}^+$  reduction reaction will be. Preventing recombination of photo-generated  $e^-$ – $h^+$  pairs in the MoS<sub>2</sub>/TNAs and g-C<sub>3</sub>N<sub>4</sub>/TNAs heterojunction structures has also been shown to increase the efficiency of PEC water splitting.

## Conclusion

MoS<sub>2</sub>/TNAs and g-C<sub>3</sub>N<sub>4</sub>/TNAs heterojunctions have been successfully fabricated for PEC water splitting. The role of g-C<sub>3</sub>N<sub>4</sub> and MoS<sub>2</sub> in mitigating the recombination of  $e^-$ – $h^+$  pairs has been demonstrated. The ability to enhance the optical interaction of the heterostructures was presented through the reduction of the bandgap energy. The outstanding application performance of g-C<sub>3</sub>N<sub>4</sub>/TNAs and MoS<sub>2</sub>/TNAs combinations was presented. In detail, the excellent water-splitting ability of g-C<sub>3</sub>N<sub>4</sub>/TNAs and MoS<sub>2</sub>/TNAs heterojunctions achieved about 139.6 and 210.6  $\mu\text{A}/\text{cm}^2$ . In addition, the PEC reaction rate was



**Figure 7:** Proposed band diagram of MoS<sub>2</sub>/TNAs (a) and g-C<sub>3</sub>N<sub>4</sub>/TNAs (b).



evaluated by the Tafel slope value, indicating a faster rate for the MoS<sub>2</sub>/TNAs heterojunction compared to the g-C<sub>3</sub>N<sub>4</sub>/TNAs heterojunction. Moreover, the photocurrent density of MoS<sub>2</sub>/TNAs is higher than that of g-C<sub>3</sub>N<sub>4</sub>/TNAs due to the strong optical interaction of the MoS<sub>2</sub>/TNAs heterostructure.

## Supporting Information

### Supporting Information File 1

Additional figures.

[<https://www.beilstein-journals.org/bjnano/content/supplementary/2190-4286-13-127-S1.pdf>]

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