



Supporting Information

for

Interfacial charge transfer processes in 2D and 3D semiconducting hybrid perovskites: azobenzene as photoswitchable ligand

Nicole Fillafer, Tobias Seewald, Lukas Schmidt-Mende and Sebastian Polarz

Beilstein J. Nanotechnol. **2020**, *11*, 466–479. doi:10.3762/bjnano.11.38

Additional experimental data

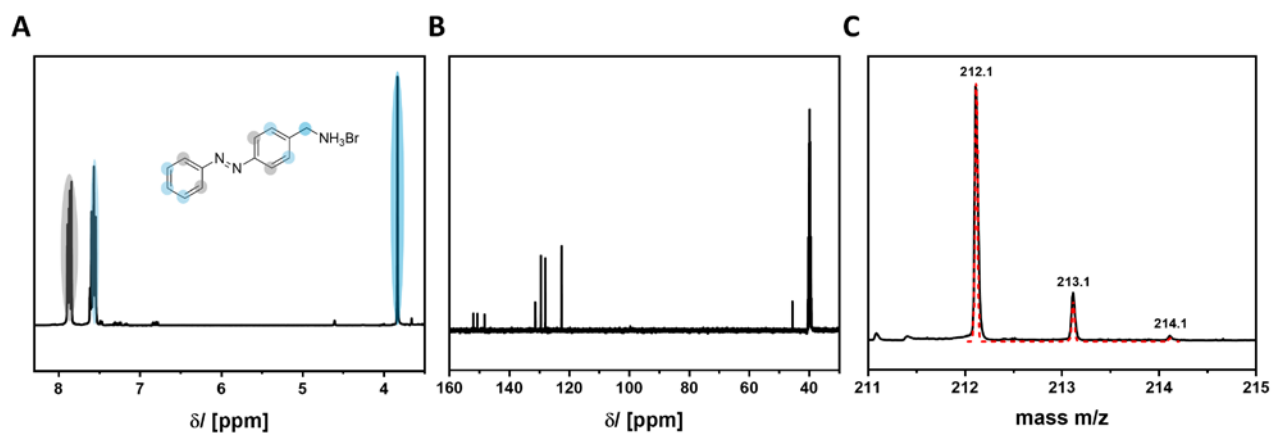


Figure S1: (A) ¹H NMR spectrum of AzoC₁ in DMSO-*d*₆, H atoms are marked. (B) ¹³C NMR spectrum of AzoC₁ in DMSO-*d*₆, (C) ESI-MS spectrum of AzoC₁, black line: experimental, red dashed line: calculated.

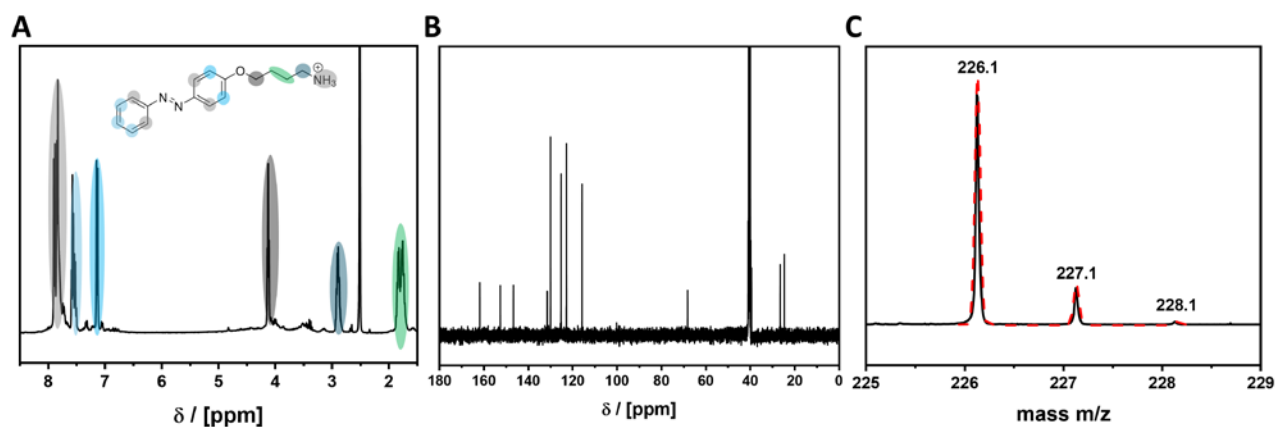


Figure S2: (A) ¹H NMR spectrum of AzoC₂ in DMSO-*d*₆, H atoms are marked. (B) ¹³C NMR spectrum of AzoC₂ in DMSO-*d*₆, (C) ESI-MS spectrum of AzoC₂, black line: experimental, red dashed line: calculated.

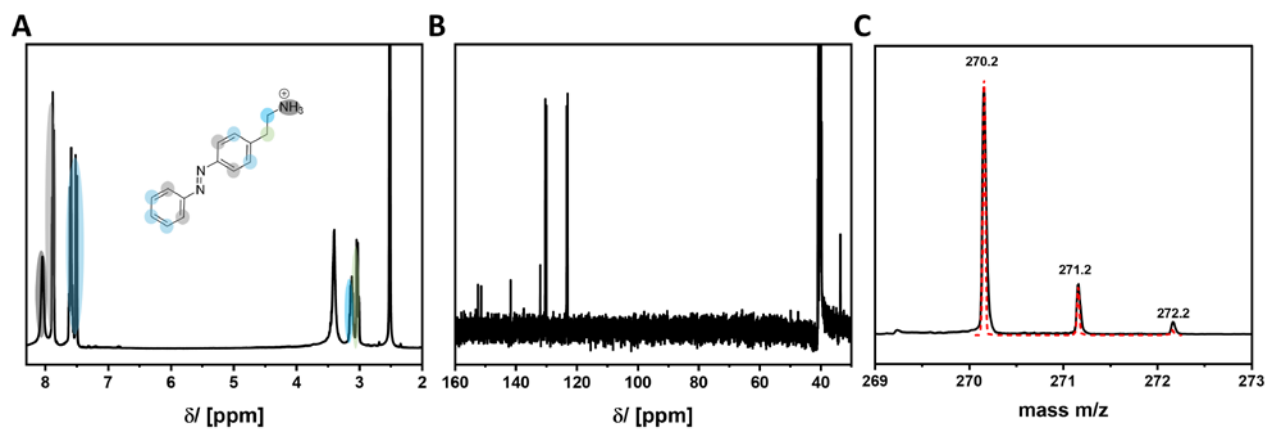


Figure S3: (A) ^1H NMR spectrum of AzoOC₄ in DMSO-*d*₆, H atoms are marked. (B) ^{13}C NMR spectrum of AzoOC₄ in DMSO-*d*₆, (C) ESI-MS spectrum of AzoOC₄, black line: experimental, red dashed line: calculated.

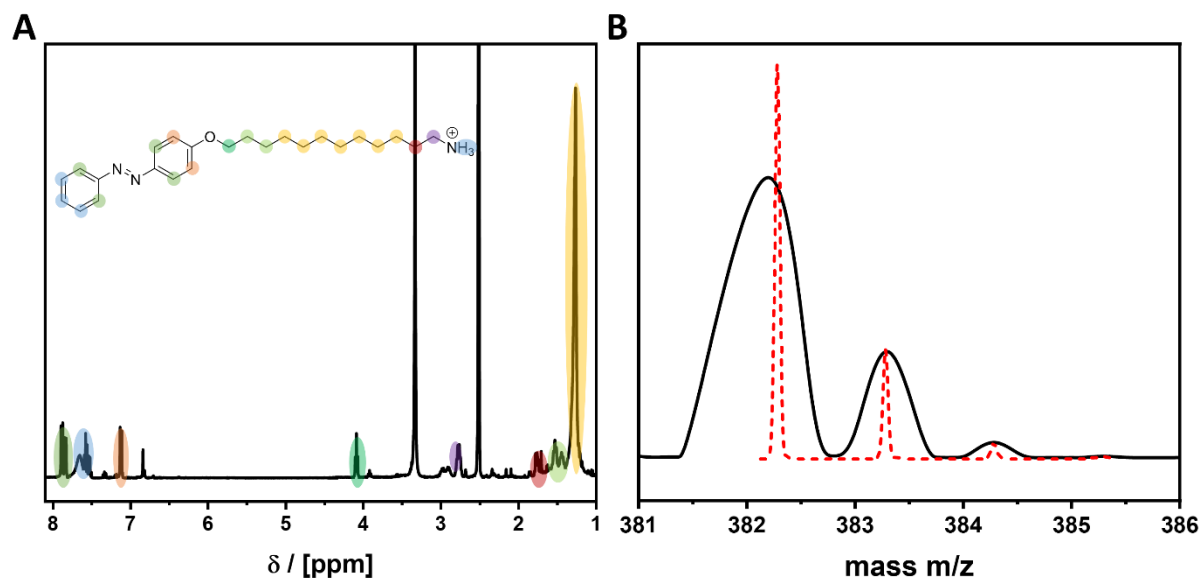


Figure S4: (A) ^1H NMR spectrum of AzoOC₁₂ in DMSO-*d*₆, H atoms are marked. ^{13}C NMR was measured but the concentration was too low. (B) ESI-MS spectrum of AzoOC₁₂, black line: experimental, red dashed line: calculated.

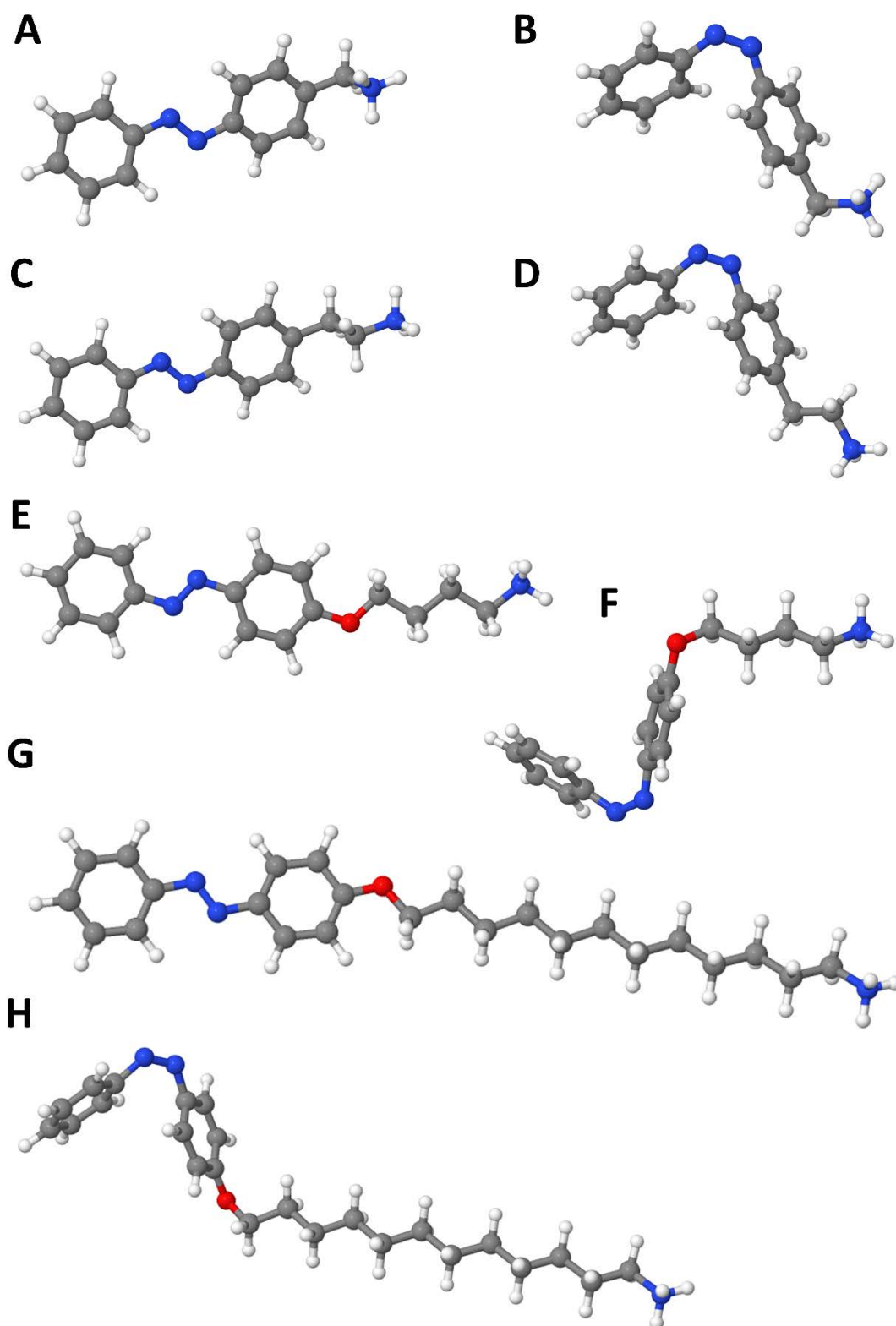


Figure S5: (A, B) DFT calculated structure of *trans*-AzoC₁ and *cis*-AzoC₁, (C, D) DFT-calculated structure of *trans*-AzoC₂ and *cis*-AzoC₂, (E, F) DFT calculated structure of *trans*-AzoOC₄ and *cis*-AzoOC₄, (G, H) DFT calculated structure of *trans*-AzoOC₁₂ and *cis*-AzoOC₁₂, all structures were created and measured with Jmol.

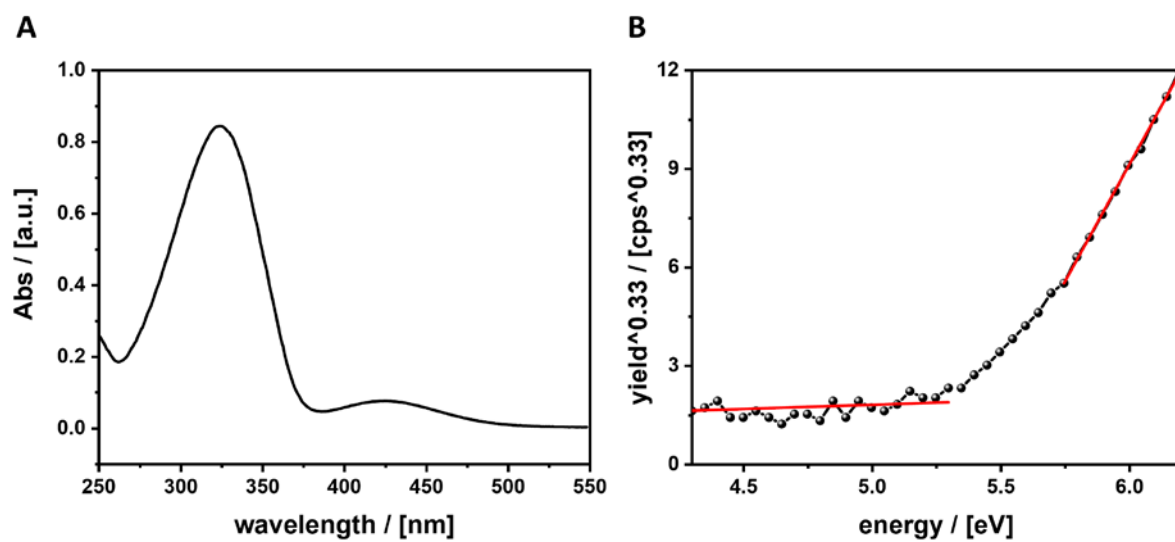


Figure S6: (A) UV-vis spectrum of AzoC₁ in H₂O, (B) PESA spectrum of AzoC₁.

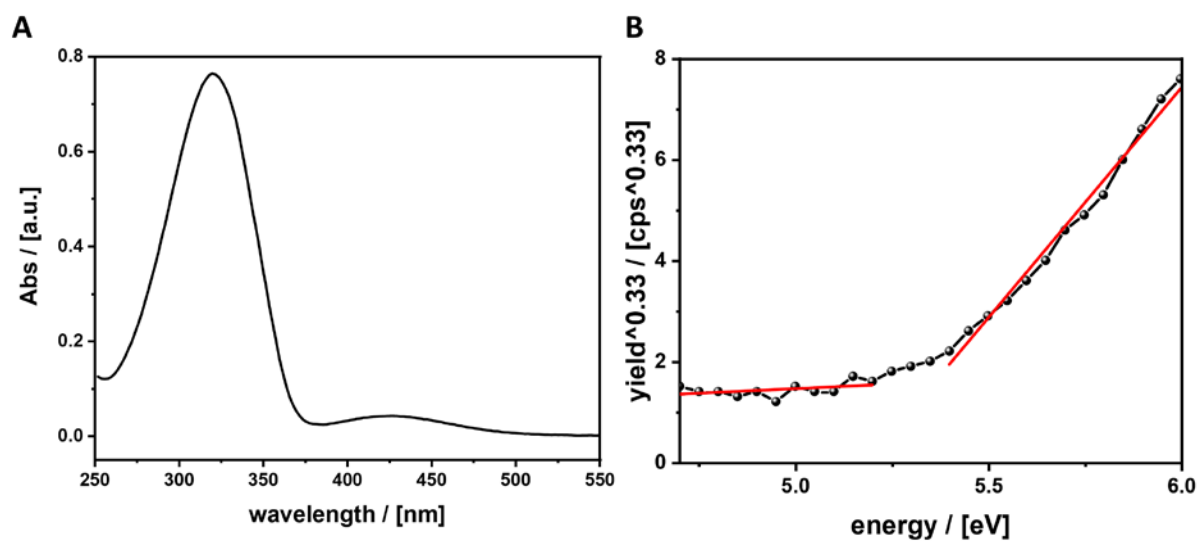


Figure S7: (A) UV-vis spectrum of AzoC₂ in H₂O, (B) PESA spectrum of AzoC₂.

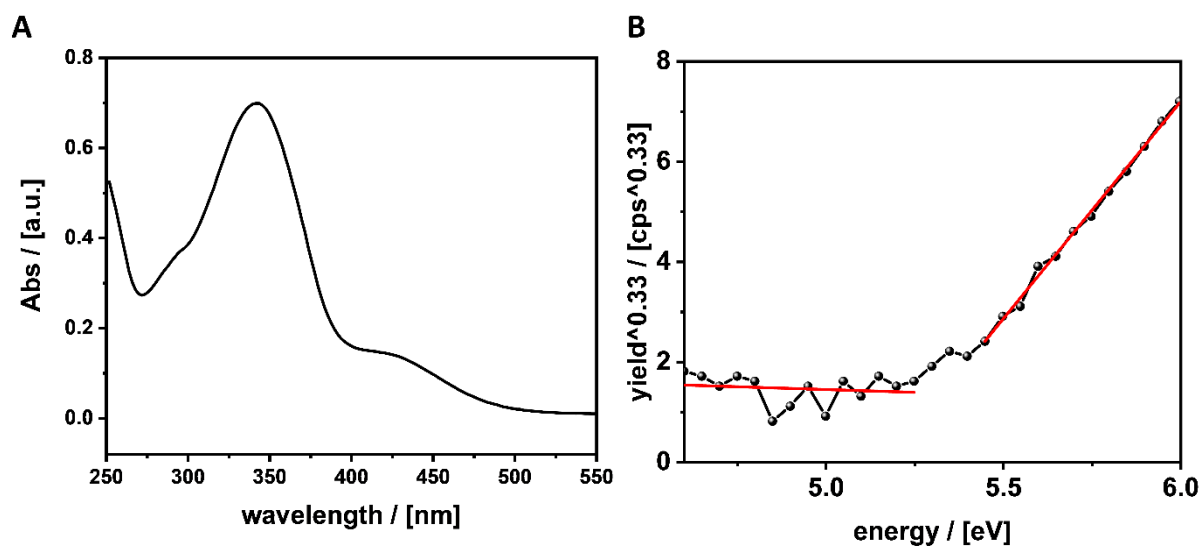


Figure S8: (A) UV-vis spectrum of AzoOC₄ in H₂O, (B) PESA spectrum of AzoOC₄.

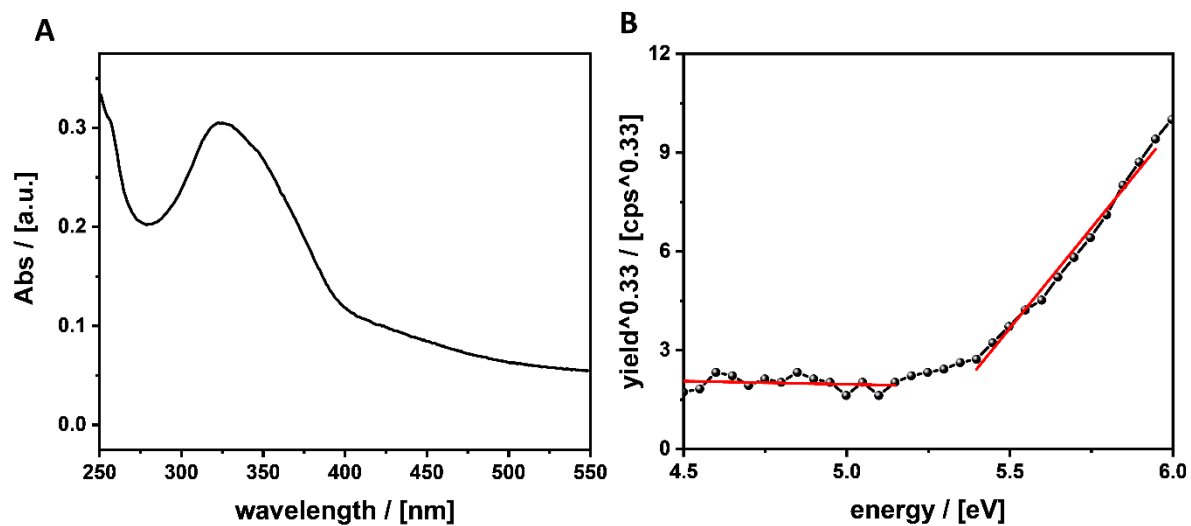


Figure S9: (A) UV-vis spectrum of AzoOC₁₂ in H₂O, (B) PESA spectrum of AzoOC₁₂.

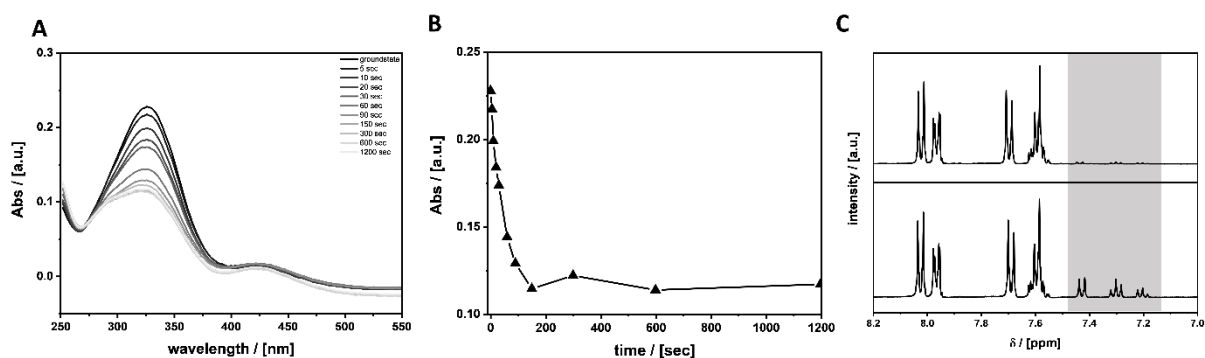


Figure S10: (A) UV–vis kinetic measurement of AzoC₁ in purified H₂O. The solution was irradiated at 313 nm, black: beginning of the reaction, light grey: ending of the reaction (B) absorption at 326 nm over a time period of 1200 s, (C) ¹H NMR of AzoC₁ in MeOD from 8.0 to 7.2 ppm, before (top) and after irradiation (bottom) at 313 nm for 1 h. Highlighted in grey is the signal of the *cis*-isomer. The DOI is 21.7%.

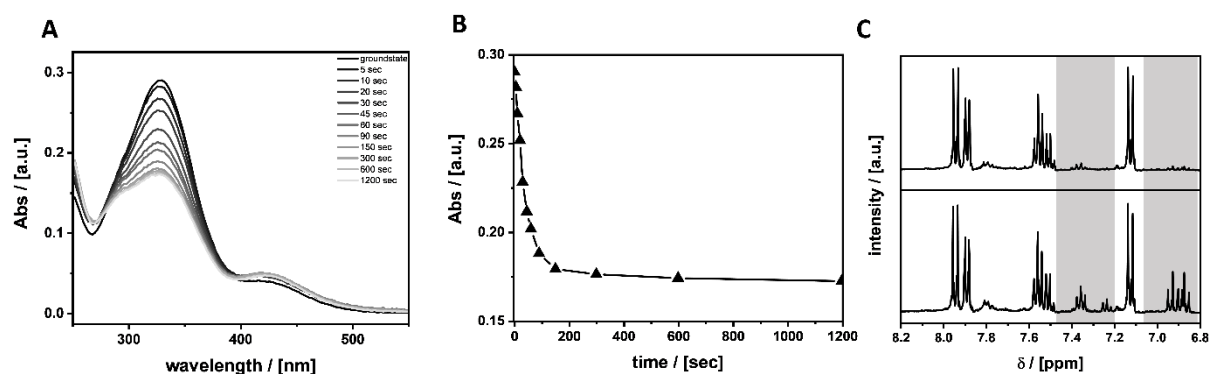


Figure S11: (A) UV–vis kinetic measurement of AzoOC₄ in purified H₂O. The solution was irradiated at 313 nm, black: beginning of the reaction, light grey: ending of the reaction (B) absorption at 328 nm over a time period of 1200 s. (C) ¹H NMR of AzoC₁ in MeOD from 8.0 to 7.2 ppm, before (top) and after irradiation (bottom) at 313 nm for 1 h. Highlighted in grey is the signal of the *cis*-isomer. The DOI is 25.1%.

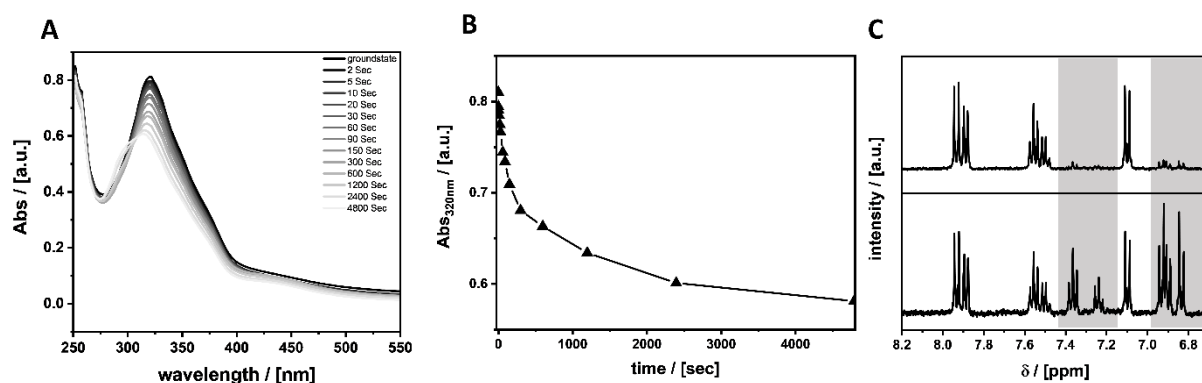


Figure S12: (A) UV–vis kinetic measurement of AzoOC₁₂ in purified H₂O. The solution was irradiated with 313 nm, black: beginning of the reaction, light grey: ending of the reaction (B) absorption at 328 nm over a time period of 4800 s. (C) ¹H NMR of AzoC₁₂ in MeOD from 8.0 to 7.2 ppm, before (top) and after irradiation (bottom) at 313 nm for 1 h. Highlighted in grey is the signal of the *cis*-isomer. The DOI is 50.7%.

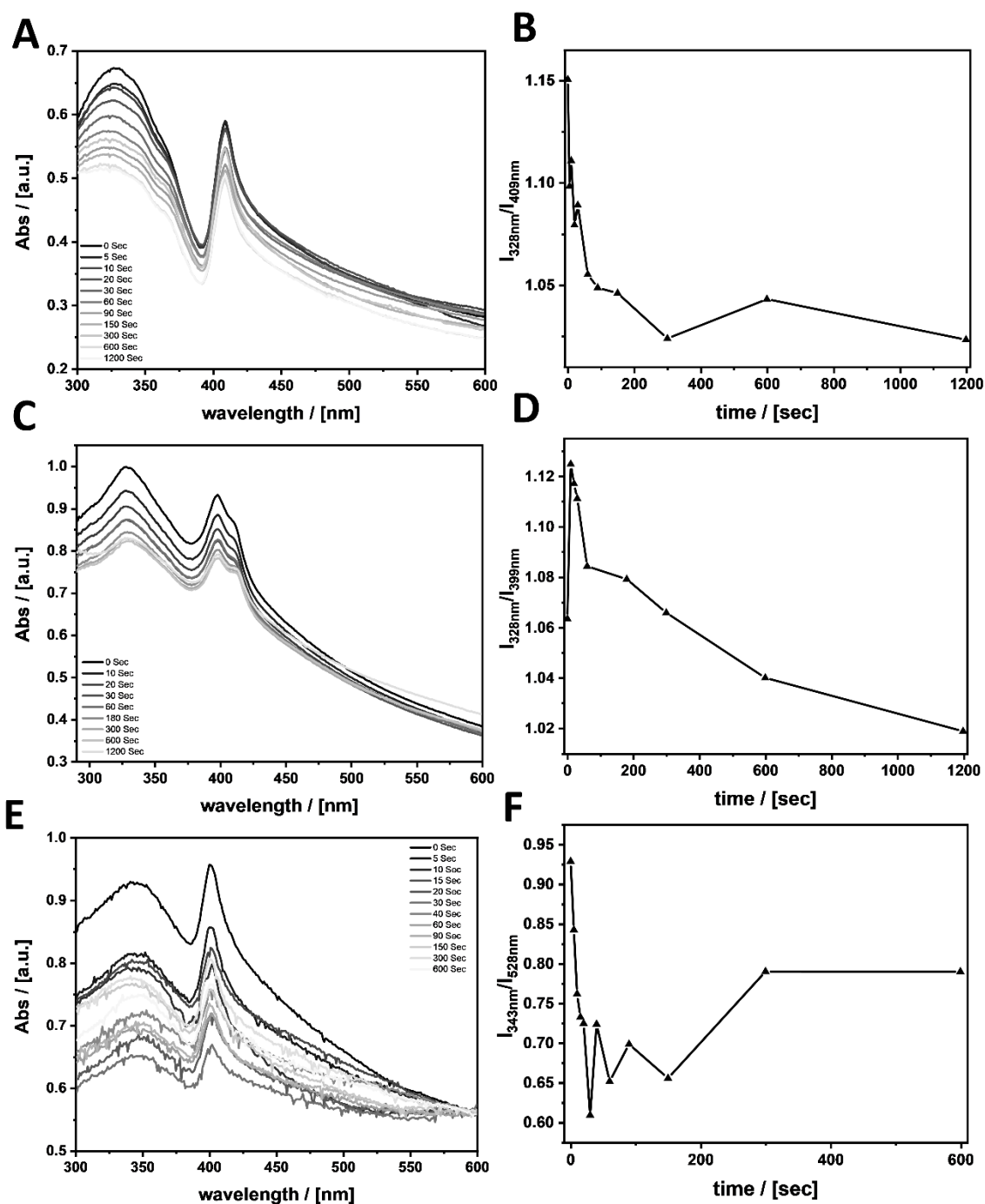


Figure S13: (A, C, E) UV–Vis measurements of dispersions of 2D-AzoC₂, 2D-AzoOC₄ and 2D-AzoOC₁₂ in toluene during irradiation at 313 nm. (B, D, F) Intensity of absorbance at the absorption maximum of the chromophore divided by the intensity of the absorbance of the excitonic bandgap signal. As the dispersions were not stable they were referenced to the signal of the bandgap.

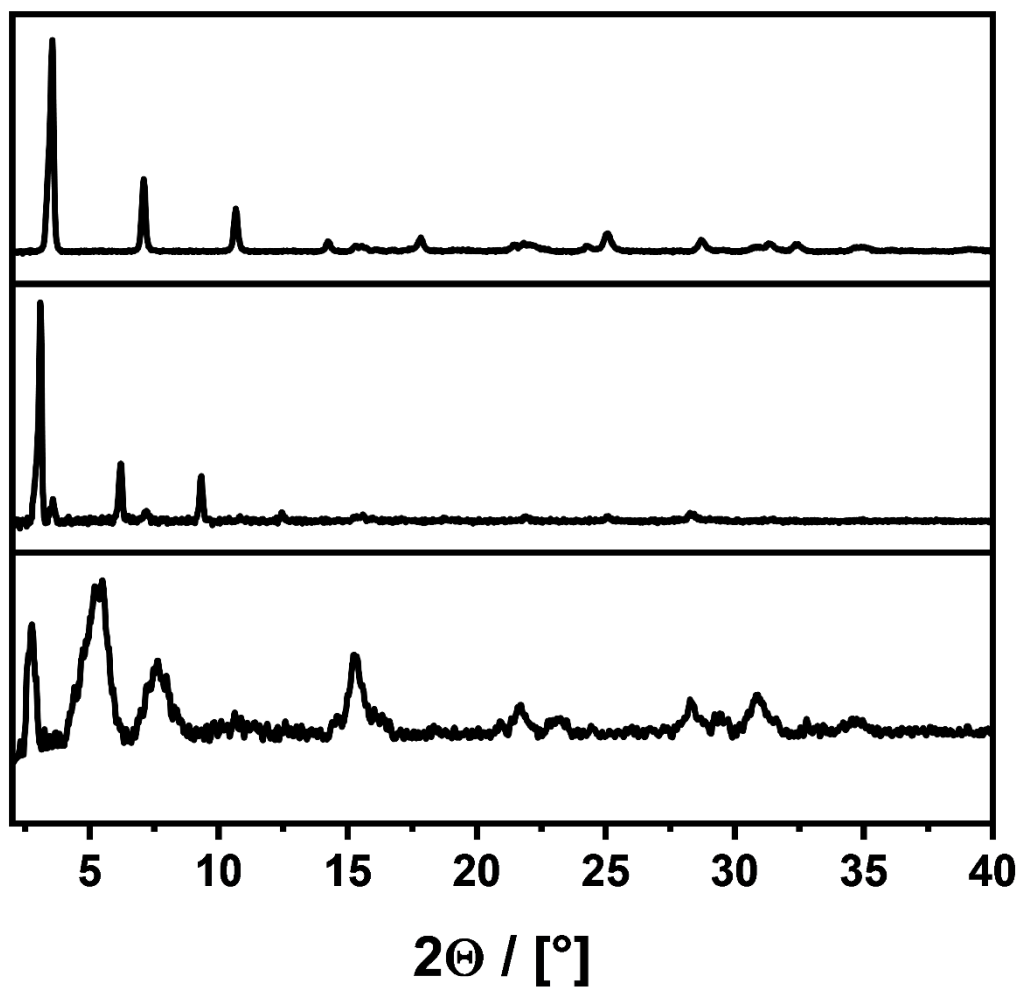


Figure S14: PXRD patterns of 2D-AzoC₂, 2D-AzoOC₄ and 2D-AzoOC₁₂ (top to bottom) after irradiation at 313 nm for 20 min. A layer spacing of $d_{(001)} = 2.48$ nm, $d_{(001)} = 2.84$ nm and $d_{(001)} = 3.40$ nm, respectively, from top to bottom was found. No change compared to Figure 2A in the main manuscript can be found.

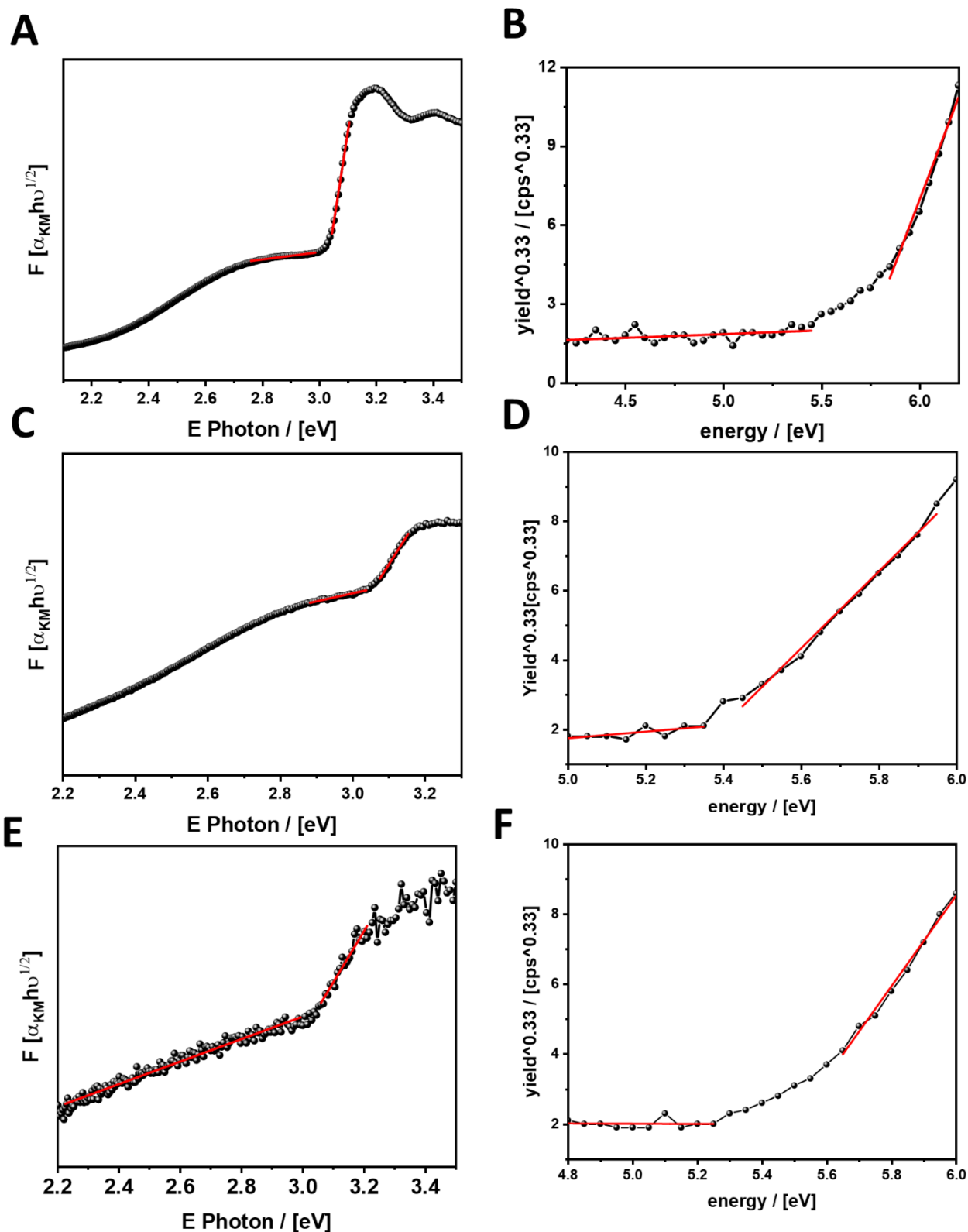


Figure S15: (A, B) Kubelka–Munk evaluation and PESA of 2D-AzoC₂, (C, D) Kubelka–Munk evaluation and PESA of 2D-AzoOC₄. (E, F) Kubelka–Munk evaluation and PESA of 2D-AzoOC₁₂. The intersection of baseline fit and slope fit gives the bandgap energy E_g or the relative energy E_{VB} of the valence band, respectively.

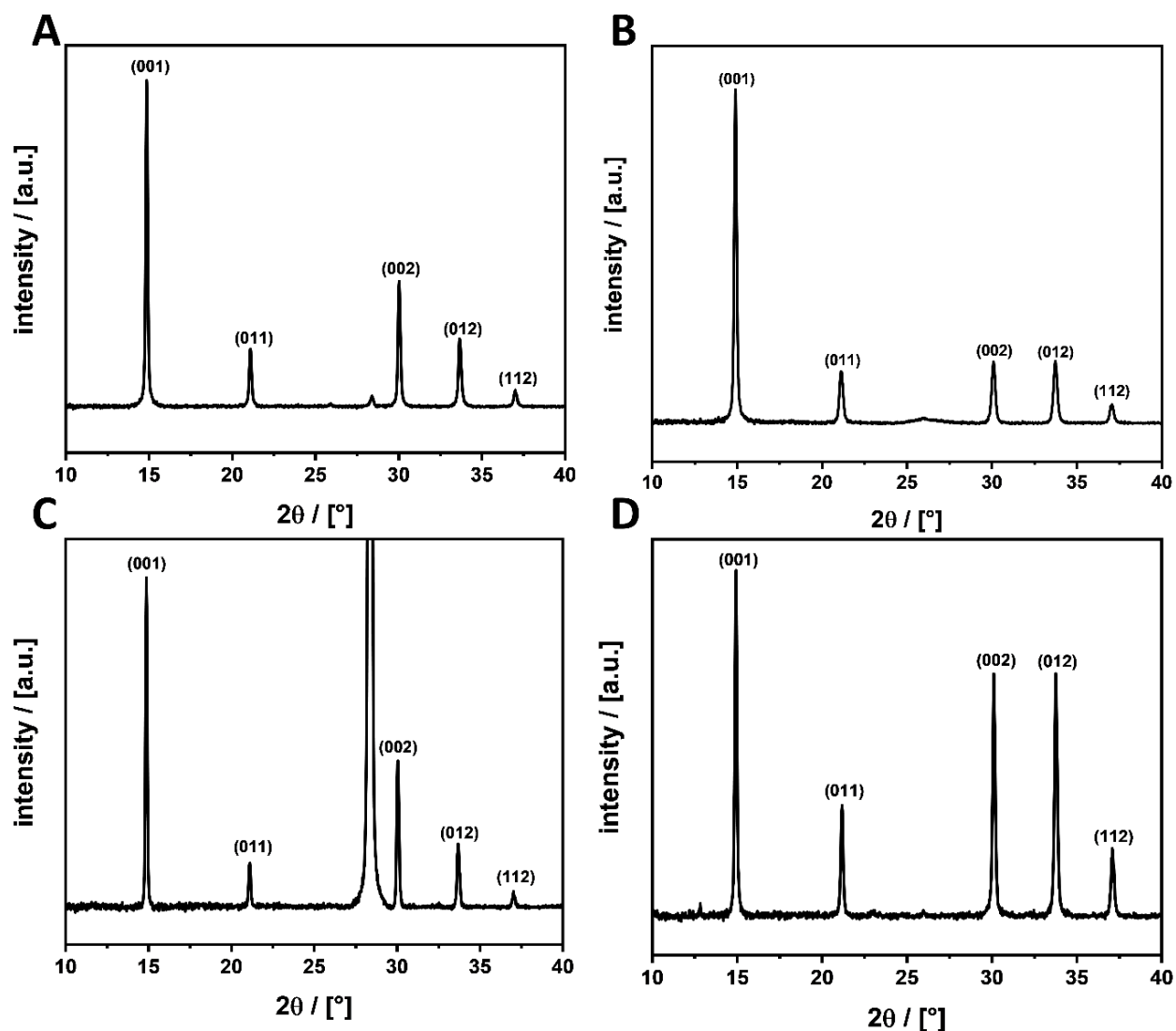


Figure S16: (A) PXR D pattern of 3D-AzoC₁, (B) PXR D pattern of 3D-AzoC₂, (C) PXR D pattern of 3D-AzoOC₄, (D) PXR D pattern of 3D-AzoOC₁₂. All reflexes originate from $\text{CH}_3\text{NH}_3\text{PbBr}_3$, the reflex at $2\theta = 28^\circ$ originates from the Si substrate.

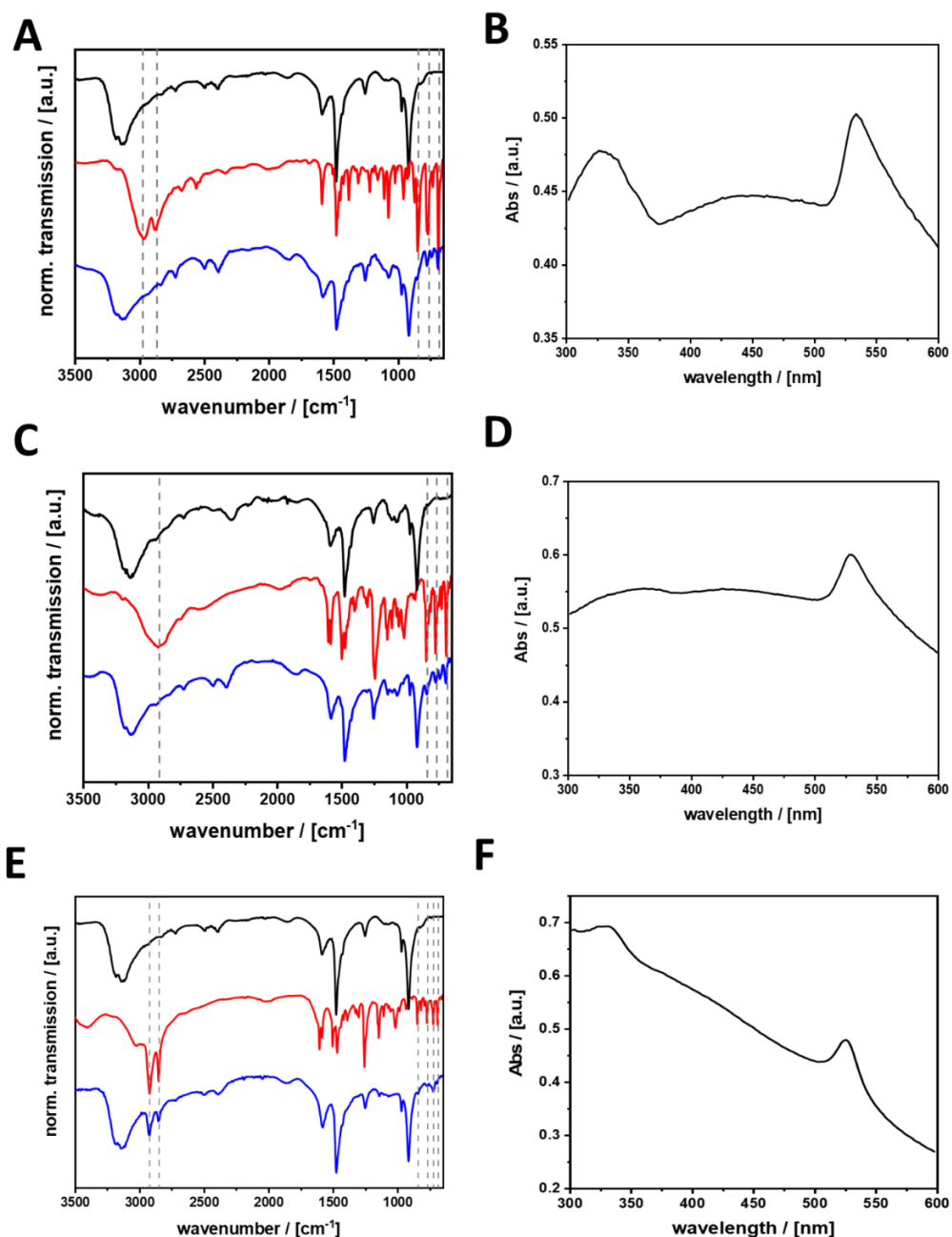


Figure S17: (A) IR spectrum of AzoC₁ (red), CH₃NH₃PbBr₃ (grey) and 3D-AzoC₁. (B) UV-vis absorption spectrum of 3D-AzoC₁ in toluene. (C) IR spectrum of AzoOC₄ (red), CH₃NH₃PbBr₃ (grey) and 3D-AzoOC₄. (D) UV-vis absorption spectrum of 3D-AzoOC₄ in toluene. (E) IR spectrum of AzoOC₁₂ (red), CH₃NH₃PbBr₃ (grey) and 3D-AzoOC₁₂. (F) UV-vis absorption spectrum of 3D-AzoOC₁₂ in toluene.

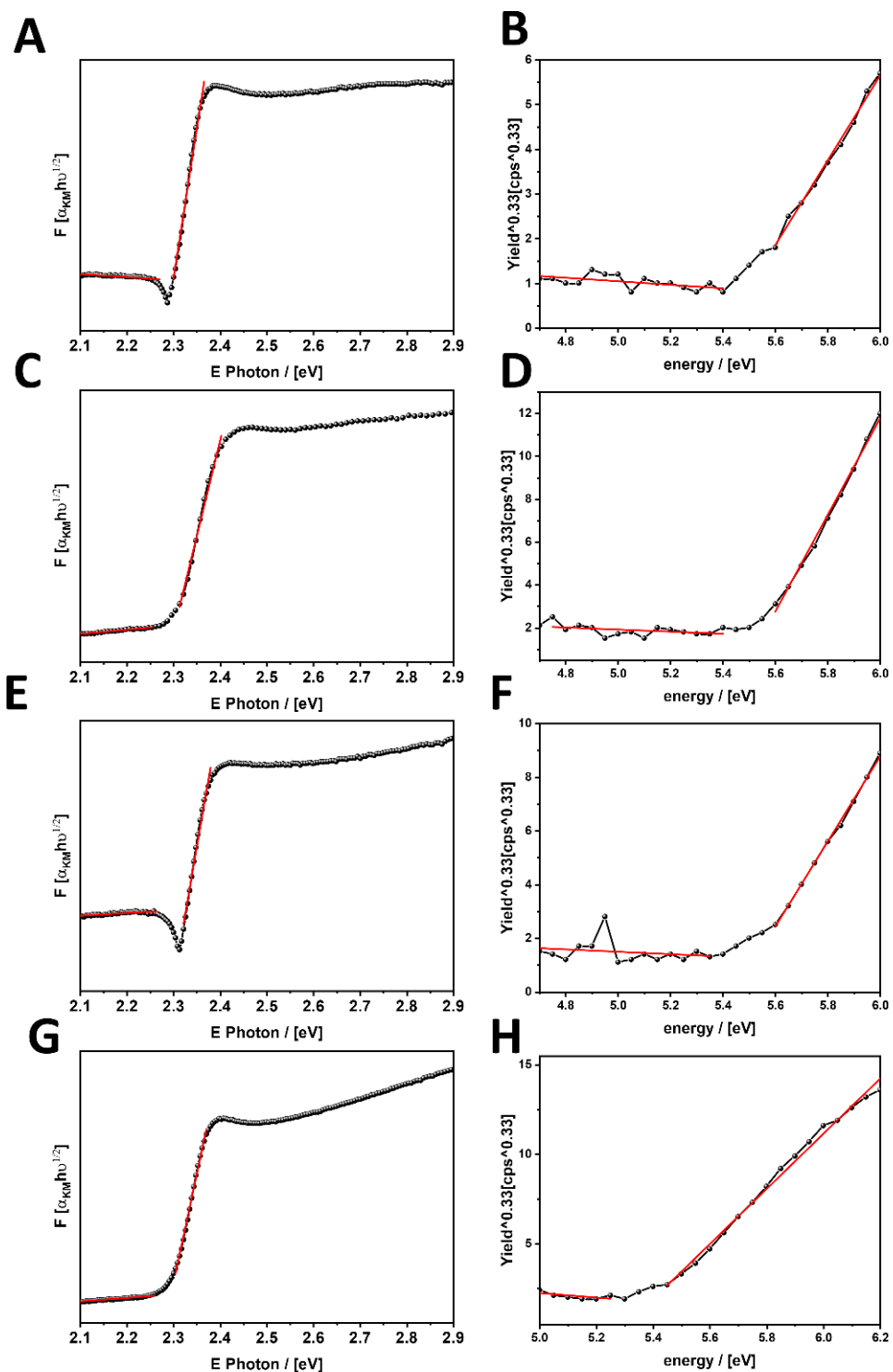


Figure S18: (A, B) UV–vis direct reflection and PESA measurement of 3D-AzoC₁. (C, D) UV–vis direct reflection and PESA measurement of 3D-AzoC₂. (E, F) UV–vis direct reflection and PESA measurement of 3D-AzoOC₄, (G, H) UV–vis direct reflection and PESA measurement of 3D-AzoOC₁₂. The intersection of baseline fit and slope fit gives the bandgap energy E_g or the relative energy E_{VB} of the valence band, respectively.

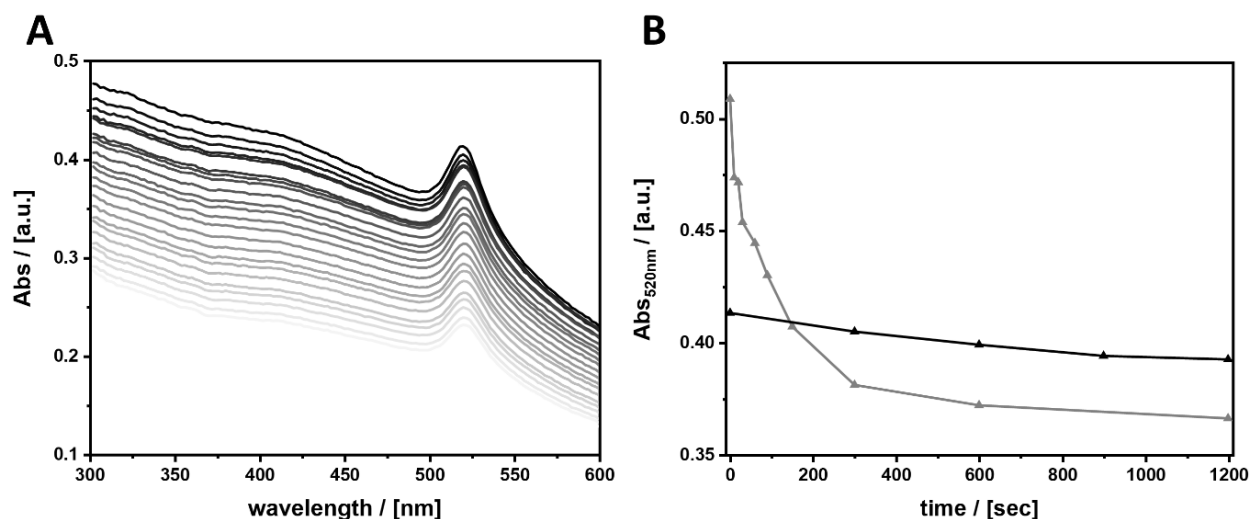


Figure S19: (A) UV-vis absorption spectra of a particle dispersion without irradiation recorded every 300 s. (B) Absorbance as a function of the time of dispersed particles without irradiation (black) and with irradiation (grey, spectrum see Figure 5 in the main manuscript). The dispersion is not stable over time, which is why we refer the azobenzene signal to the signal of the bandgap.

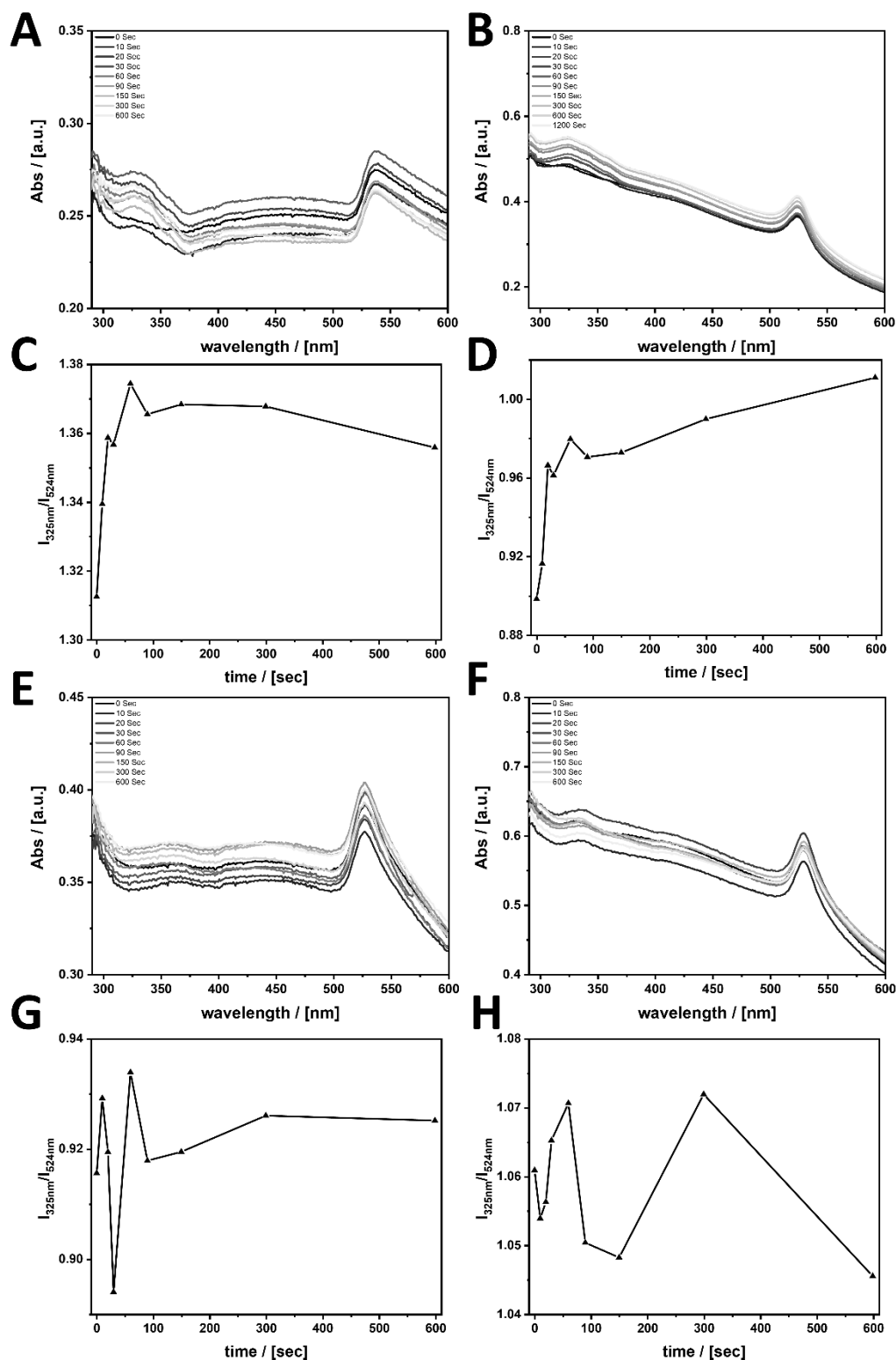


Figure S20: (A, B, E, F) UV–vis spectra of dispersions (irradiated for 5 min at 313 nm) of 3D-AzoC₁, 3D-AzoC₂, 3D-AzoOC₄ and 3D-AzoOC₁₂ in toluene while the particles are irradiated at 438 nm. (C, D, G, H) Intensity of absorbance at maximum of AzoC₁, AzoC₂, AzoOC₄ and AzoOC₁₂ divided by the intensity of the absorbance of the excitonic bandgap signal. As the dispersions were not stable they were referenced to the signal of the bandgap.

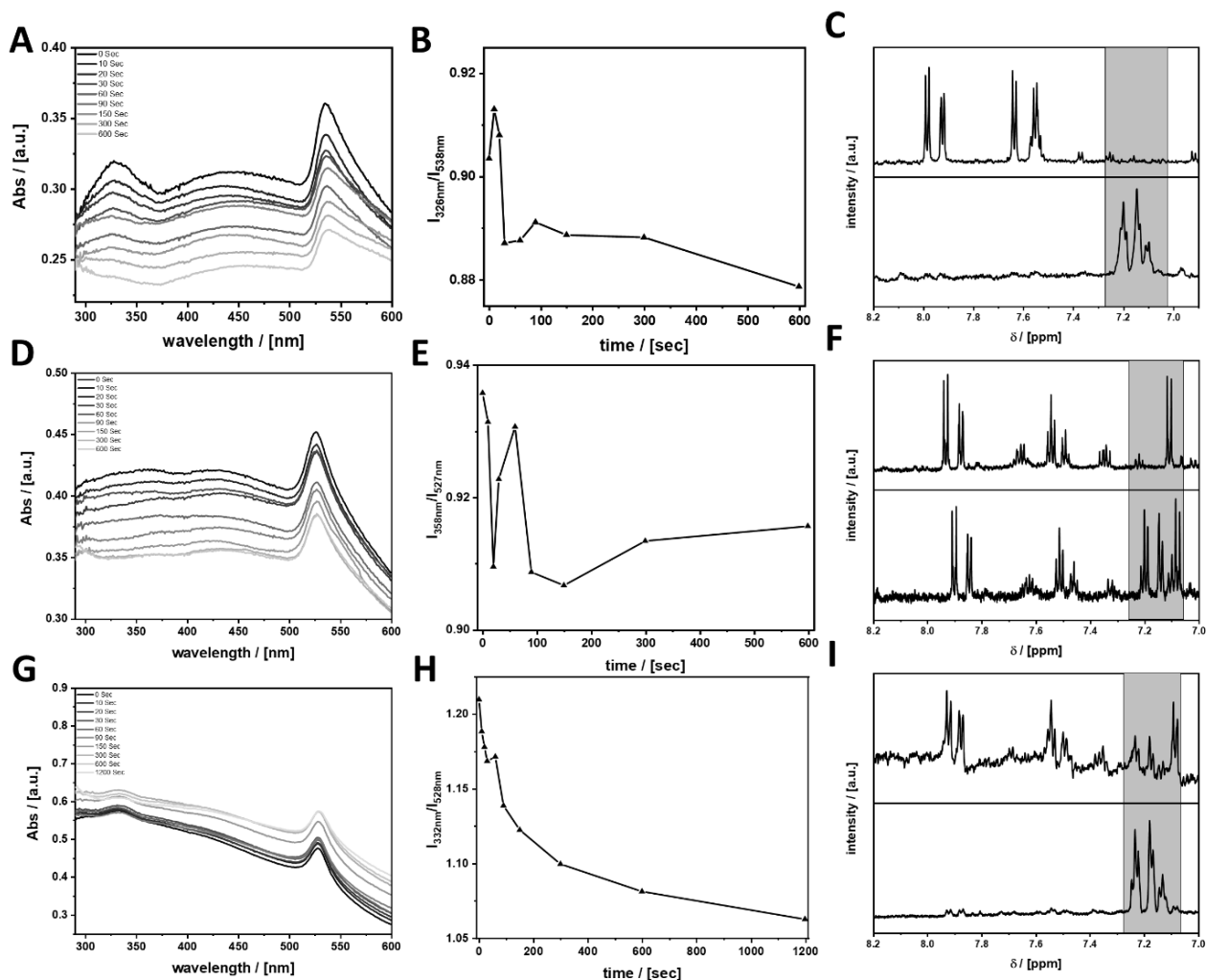


Figure S21: (A, B) UV–vis kinetic measurement of 3D-AzoC₁ in toluene over a period of time of 600 s. Because the dispersion is not stable, the absorbance of AzoC₁ at 326 nm is referenced to the absorbance of the excitonic bandgap at 538 nm and plotted as a function of the time. (C) ^1H NMR spectrum of dissolved 3D-AzoC₁ before (top) and after irradiation at 313 nm for 600 s (bottom). (D, E) UV–Vis kinetic measurement of 3D-AzoOC₄ in toluene over a period of time of 600 s. Because the dispersion is not stable, the absorbance of AzoOC₄ at 358 nm is referenced to the absorbance of the excitonic bandgap at 527 nm and plotted as a function of the time. (F) ^1H NMR spectrum of dissolved 3D-AzoOC₄ before (top) and after irradiation at 313 nm for 600 s (bottom). (G, H) UV–vis kinetic measurement of 3D-AzoOC₁₂ in toluene over a period of time of 1200 s. Because the dispersion is not stable, the absorbance of AzoOC₁₂ at 332 nm is referenced to the absorbance of the excitonic bandgap at 528 nm and plotted as a function of the time. (I) ^1H NMR spectrum of dissolved 3D-AzoOC₁₂ before (top) and after irradiation at 313 nm for 1200 s (bottom).

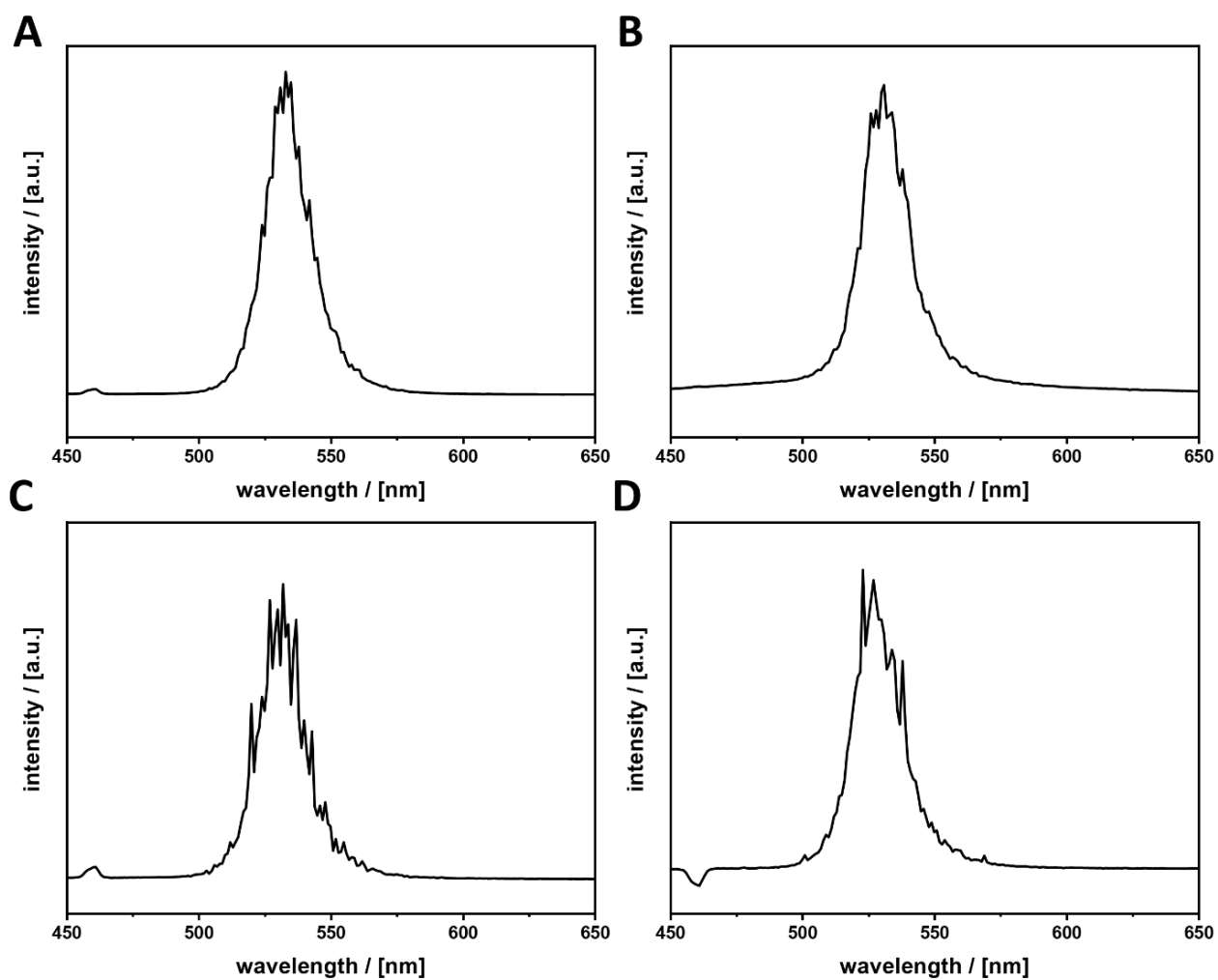


Figure S22: PL spectra of (A) 3D-AzoC₁, (B) 3D-AzoC₂, (C) AzoOC₄ and (D) AzoOC₁₂ dispersed in toluene, with $\lambda_{\text{exc}} = 405$ nm. To filter out the excitation wavelength a cut-off filter at 420 nm was used. Signals at 460 nm originate from toluene.

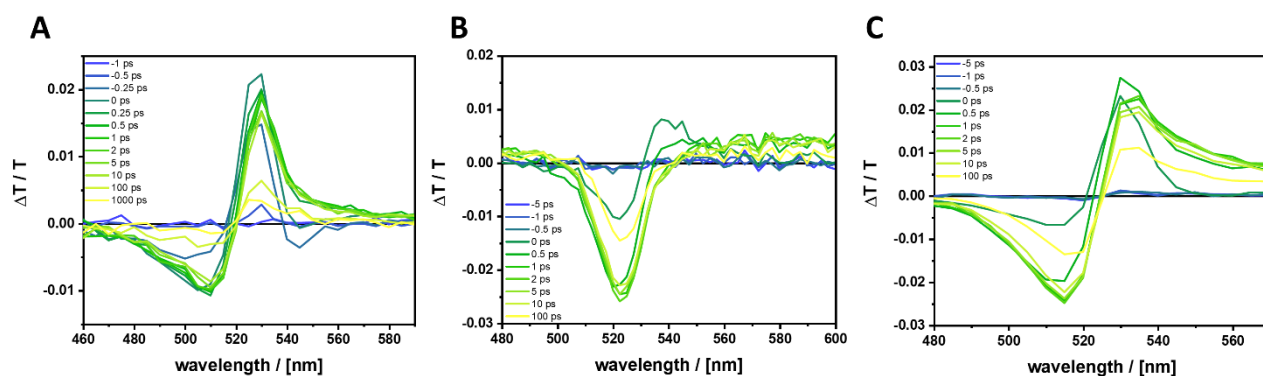


Figure S23: Transient absorption spectra of (A) 3D-AzoC₁, (B) 3D-AzoC₂ and (C) 3D-AzoOC₁₂. All samples were excited with a 398 nm pulse. Negative signals are obtained due to refractive index changes, which are caused by a change of the bandgap energy. Positive signals show a ground-state bleach, as the absorption at the bandgap decreases due to excited states.